Nonpositively curved metric in the positive cone of a finite von Neumann algebra*

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Abstract

In this paper we study the metric geometry of the space Σ of positive invertible elements of a von Neumann algebra \mathcal{A} with a finite, normal and faithful tracial state τ . The trace induces an incomplete Riemannian metric $\langle x,y \rangle_a = \tau(ya^{-1}xa^{-1})$, and though the techniques involved are quite different, the situation here resembles in many relevant aspects that of the $n \times n$ matrices when they are regarded as a symmetric space. For instance we prove that geodesics are the shortest paths for the metric induced, and that the geodesic distance is a convex function; we give an intrinsic (algebraic) characterization of the geodesically convex submanifolds M of Σ , and under suitable hypothesis we prove a factorization theorem for elements in the algebra that resembles the Iwasawa decomposition for matrices. This factorization is obtained via a nonlinear orthogonal projection $\Pi_M : \Sigma \to M$, a map which turns out to be contractive for the geodesic distance.

1 Introduction

Let \mathcal{A} be a von Neumann algebra with a finite (normal, faithful) trace τ . Denote by \mathcal{A}_h the set of selfadjoint elements of \mathcal{A} , by $G_{\mathcal{A}}$ the group of invertible elements, and by Σ the set

$$\Sigma = e^{\mathcal{A}_h} = \{ a \in G_{\mathcal{A}} : a \ge 0 \};$$

 Σ is an open subset of \mathcal{A}_h in the norm topology. Therefore if one regards it as a manifold, its tangent spaces identify with \mathcal{A}_h . We endow these tangent spaces with the (incomplete) Hilbert-Riemann metric

$$< x, y>_a = \tau(xa^{-1}ya^{-1}), \quad a \in \Sigma, \ x, y \in \mathcal{A}_h.$$
 (1)

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Note that $||x||_a^2 = \langle x, x \rangle_a := \tau(xa^{-1}xa^{-1})$, and also that this metric is invariant for the action $I_g: x \mapsto gxg^*$, where $g \in G_A$.

As in classical differential geometry, one obtains a metric d for Σ by considering

$$dist(a,b) = \inf\{Length(\gamma) : \gamma \text{ is a smooth curve joining } a \text{ and } b\},$$
 (2)

where smooth means differentiable in the norm induced topology and the length of a curve $\gamma(t)$, $t \in [0, 1]$ is measured using the inner product above (1):

$$Length(\gamma) = \int_0^1 \langle \dot{\gamma}(t), \dot{\gamma}(t) \rangle_{\gamma(t)}^{\frac{1}{2}} dt.$$

The purpose of this paper is the geometric study of the resulting metric space, and particularly, of its convex subsets.

If \mathcal{A} is finite dimensional, i.e. a sum of matrix spaces, this metric is well known: it is the non positively curved Riemannian metric on the set of positive definite matrices, which is a universal model space for (finite dimensional) non positively curved manifolds on non compact type (see [4] and [5]).

If \mathcal{A} is of type II₁, the trace inner product is not complete, so that Σ , with the inner products \langle , \rangle_a , is not a Hilbert-Riemann manifold properly speaking. For instance, the exponential map

$$exp: \mathcal{A}_h \to \Sigma, \ exp(x) = e^x,$$

which is a global diffeomorphism in the norm topology, is continuous but non differentiable in the 2-norm $\| \|_2$ induced by τ (namely $\|x\|_2 = \tau(x^*x)^{\frac{1}{2}}$). The set Σ itself is not a differentiable manifold with this norm.

However, the metric space $(\Sigma, dist)$ behaves in many senses like in the finite dimensional setting. Let us mention a few issues:

- 1. The isometric action of the group G_A via $g \mapsto I_q$, where $I_q(x) = gxg^*$
- 2. Minimality of geodesics (i.e. solutions of Euler's equation are minimizing for the distance introduced above in (2), see Theorem 3.1)
- 3. Convexity of the map $t \mapsto dist(\gamma(t), \delta(t))$ which gives distance among geodesics (Corollary 3.4)
- 4. Algebraic structure of (geodesically) convex subsets (Theorem 4.4).
- 5. Normal projections to convex submanifolds and their minimality (Lemma 5.3 and Theorem 5.4)
- 6. Existence and uniqueness of a factorization for invertible elements by means of convex submanifolds (Corollary 5.7)

2 Main inequalities

The following inequality will be useful; its proof for $n \times n$ real matrices can be found in the inspiring paper of G.D. Mostow [5]. It is called by R. Bhatia [1] the exponential metric increasing property. Bhatia proves it for matrices (and for more general norms). However his proof for the 2-norm is valid almost verbatim in the infinite dimensional context for an arbitrary (finite, faithful) tracial state. We transcribe it. We use the fact that selfadjoint elements in a von Neumann algebra can be approximated by selfadjoint elements with finite spectrum.

Lemma 2.1. Let τ be a tracial faithful state in A, and x, y selfadjoint elements of A. If exp denotes the usual exponential map, $\exp(x) = e^x$, then

$$||y||_2 \le ||e^{-x} dexp_x(y)||_2.$$
 (3)

Proof. First we must establish the formula

$$dexp_x(y) = \int_0^1 e^{tx} y e^{(1-t)x} dt.$$

Note that $dexp_x(y) = \frac{d}{dt}e^{x+ty}|_{t=0}$. Then

$$dexp_x(y) = y + \frac{1}{2}(yx + xy) + \frac{1}{6}(yx^2 + xyx + x^2y) + \dots$$

On the other hand,

$$e^{tx}ye^{(1-t)x} = y + txy + (1-t)yx + \frac{1}{2}(1-t)^2yx^2 + t(1-t)xyx + \frac{1}{2}t^2x^2y + \dots$$

Integrating this series (which is absolutely convergent) term by term proves the equality. Denote $a = e^x$. Let us show now that if b is positive in A,

$$||a^{\frac{1}{2}}ba^{\frac{1}{2}}||_{2} \le ||\int_{0}^{1} a^{t}ba^{1-t}dt||_{2}.$$

$$\tag{4}$$

Assume first that a has finite spectrum: $a = \sum_{i=1}^{n} \alpha_i p_i$, with $\alpha_i > 0$ and $\sum_{i=1}^{n} p_i = 1$. Then $a^{\frac{1}{2}}ba^{\frac{1}{2}} = \sum_{i,j=1}^{n} \alpha_i^{\frac{1}{2}} \alpha_j^{\frac{1}{2}} p_i b p_j$. Therefore

$$\|a^{\frac{1}{2}}ba^{\frac{1}{2}}\|_{2}^{2} = \sum_{i,j=1}^{n} \alpha_{i}\alpha_{j}\tau(p_{i}bp_{j}bp_{i}).$$

Analogously $\int_0^1 a^t b a^{1-t} dt = \sum_{i,j}^n \int_0^1 \alpha_i^t \alpha_j^{1-t} dt \ p_i b p_j$ and

$$\|\int_0^1 a^t b a^{1-t} dt\|_2^2 = \sum_{i,j}^n \int_0^1 \alpha_i^{2t} \alpha_j^{2(1-t)} dt \ \tau(p_i b p_j b p_i) = \sum_{i,j}^n \frac{\alpha_i^2 - \alpha_j^2}{2 \ln \alpha_i - 2 \ln \alpha_j} \tau(p_i b p_j b p_i).$$

Note that $p_i b p_j b p_i$ is positive. Also one has the elementary inequality

$$\sqrt{st} \le \frac{s-t}{\ln s - \ln t}$$

for s, t > 0. Then

$$\alpha_i \alpha_j p_i b p_j b p_i \le \frac{\alpha_i^2 - \alpha_j^2}{2 \ln \alpha_i - 2 \ln \alpha_j} p_i b p_j b p_i.$$

Taking traces and adding yields (4) in this case. In the general case, the inequality follows by approximating (in norm) the element a with a positive elements with finite spectrum.

As in [1], put $b = e^{-x/2}ye^{-x/2}$ in (4):

$$||y||_{2} \le ||\int_{0}^{1} e^{tx} (e^{-x/2} y e^{-x/2}) e^{(1-t)x} dt||_{2} = ||e^{-x/2} \int_{0}^{1} e^{tx} y e^{(1-t)x} dt e^{-x/2}||_{2}$$
$$= ||e^{-x/2} (dex p_{x}(y)) e^{-x/2}||_{2}.$$

If a is positive and invertible and b is selfadjoint, by the Cauchy-Schwarz inequality for τ , one has

$$\|a^{-\frac{1}{2}}ba^{-\frac{1}{2}}\|_2^2 = \tau(a^{-1}ba^{-1}b) \le \tau(a^{-1}b^2a^{-1})^{\frac{1}{2}}\tau(ba^{-2}b)^{\frac{1}{2}} = \|a^{-1}b\|_2^2.$$

Using this inequality for $a = e^x$ and $b = dexp_x(y)$ one obtains

$$||y||_2 \le ||e^{-x/2}(dexp_x(y))e^{-x/2}||_2 \le ||e^{-x}(dexp_x(y))||_2.$$

Corollary 2.2. For any $x \in A_h$, the map $T_x : y \mapsto e^{-x/2} dexp_x(y)e^{-x/2}$ is bounded, symmetric for the 2-inner product (when restricted to A_h) and invertible. The inverse is contractive i.e $T_x^{-1}(z)|_2 \le ||z||_2$.

Proof. The map is clearly bounded and invertible, the bound for the inverse follows from the proof of the previous Lemma. To prove that it is symmetric, note that

$$\langle T_x(y), z \rangle_2 = \tau(zT_x(y)) = \tau(e^{-x/2} \sum_{n \ge 0} \frac{1}{n!} \sum_{p+q=n-1} x^p y x^q e^{-x/2} z) =$$

$$= \sum_{n \ge 0} \frac{1}{n!} \sum_{p+q=n-1} \tau(e^{-x/2} x^p y x^q e^{-x/2} z) = \sum_{n \ge 0} \frac{1}{n!} \sum_{p+q=n-1} \tau(x^p e^{-x/2} y e^{-x/2} x^q z) =$$

$$= \sum_{n \ge 0} \frac{1}{n!} \sum_{p+q=n-1} \tau(e^{-x/2} x^q z x^p e^{-x/2} y) = \tau(T_x(z)y) = \langle y, T_x(z) \rangle_2 .$$

3 Geodesic distance

For X, Y smooth vector fields in Σ and $p \in \Sigma$, we introduce the expression

$$(\nabla_X Y)_p = \{X(Y)\}_p - \frac{1}{2} (X_p p^{-1} Y_p + Y_p p^{-1} X_p)$$
 (5)

where X(Y) denotes derivation of the vector field Y in the direction of X (performed in the linear space \mathcal{A}_h). Note that ∇ is clearly symmetric and verifies all the formal identities of a connection. The compatibility condition between the connection and the metric

 $\frac{d}{dt} < X, Y >_{\gamma} = < \nabla_{\dot{\gamma}} X, Y >_{\gamma} + < X, \nabla_{\dot{\gamma}} Y >_{\gamma}$

is fulfilled for any smooth curve $\gamma \subset \Sigma$ and X, Y tangent vector fields along γ . This identity is straightforward from the definitions for both terms and the cyclicity of the trace. This says that ∇ is the "Levi-Civita" connection of the metric introduced. Euler's equation $\nabla_{\dot{\gamma}}\dot{\gamma} = 0$ reads $\ddot{\gamma} = \dot{\gamma}\gamma^{-1}\dot{\gamma}$, and it is easy to see that the (unique) solution of this equation with $\gamma(0) = p, \gamma(1) = q$ is given by the curve

$$\delta_{pq}(t) = p^{\frac{1}{2}} \left(p^{-\frac{1}{2}} q p^{-\frac{1}{2}} \right)^t p^{\frac{1}{2}}. \tag{6}$$

Note that $\delta_{pq} \subset \Sigma$ because aba is positive invertible whenever a, b are positive invertible.

We will prove that the shortest path joining p to q is given by the formula above (Theorem 3.1); these curves look formally equal to the geodesics between positive definite matrices (regarded as a symmetric space).

We will use Exp_p to denote the exponential map of Σ . Note that

$$\operatorname{Exp}_p(v) = p^{\frac{1}{2}} e^{p^{-\frac{1}{2}} v p^{-\frac{1}{2}}} p^{\frac{1}{2}}.$$

Rearranging the exponential series we get a simpler expression

$$\operatorname{Exp}_{p}(v) = p e^{p^{-1}v} = e^{vp^{-1}}p.$$

A straightforward computation also shows that for $p, q \in \Sigma$ we have

$$\operatorname{Exp}_p^{-1}(q) = p^{\frac{1}{2}} \ln(p^{-\frac{1}{2}} q p^{-\frac{1}{2}}) p^{\frac{1}{2}}.$$

As mentioned in the introduction, we measure curves in Σ using the norms in the tangent space, namely

$$Length(\alpha) = \int_0^1 \|\dot{\alpha}(t)\|_{\alpha(t)} dt.$$

We have $\|\delta_{p,q}(t)\|_{\delta_{p,q}(t)} \equiv \|\ln(p^{-\frac{1}{2}}qp^{-\frac{1}{2}})\|_2$, so for the geodesics introduced in equation (6), we have $L(\delta_{p,q}) = \|\ln(p^{-\frac{1}{2}}qp^{-\frac{1}{2}})\|_2$.

Theorem 3.1. Let $a, b \in \Sigma$. Then the geodesic $\delta_{a,b}$ is the shortest curve joining a and b in Σ , if the length of curves is measured with the metric defined above.

Proof. Let γ be a smooth curve in Σ with $\gamma(0) = a$ and $\gamma(1) = b$. We must compare the length of γ with the length of $\delta_{a,b}$. Since the invertible group acts isometrically for the metric, it preserves the lengths of curves. Thus we way act with $a^{-\frac{1}{2}}$, and suppose that both curves start at 1, or equivalently, a = 1. Therefore $\delta_{1,b}(t) = \delta(t) = e^{tx}$, with $x = \ln b$. The length of δ is therefore $\tau(x^2)^{\frac{1}{2}} = ||x||_2$. The proof follows easily from the inequality proved above. Indeed, since γ is a smooth curve in Σ , it is of the form $\gamma(t) = e^{\alpha(t)}$, with $\alpha = \ln \gamma$. Then α is a smooth curve of selfadjoints with $\alpha(0) = 0$ and $\alpha(1) = x$. Moreover,

$$\tau((\gamma^{-1}\dot{\gamma}\gamma^{-1}\dot{\gamma})^{\frac{1}{2}} = \|e^{-\alpha}e^{\dot{\alpha}}\|_{2} = \|e^{-\alpha}dexp_{\alpha}(\dot{\alpha})\|_{2}.$$

By the inequality in the above lemma, this is not smaller than $\|\dot{\alpha}\|_2$. Then

$$\int_0^1 \tau ((\gamma^{-1}\dot{\gamma}\gamma^{-1}\dot{\gamma})^{\frac{1}{2}}dt \ge \int_0^1 \|\dot{\alpha}\|_2 dt \ge \|\int_0^1 \dot{\alpha} dt \|_2 = \|x\|_2 = \tau(x^2)^{\frac{1}{2}}.$$

Remark 3.2. The geodesic distance induced by the metric is given by

$$dist(a,b) = \tau \left(\ln(a^{-\frac{1}{2}}ba^{-\frac{1}{2}})^2 \right)^{\frac{1}{2}}.$$

The curvature tensor [2] is given by

$$R_a(x,y)z = -\frac{1}{4}a[[a^{-1}x, a^{-1}y], a^{-1}]$$

where [,] is the usual commutator, i.e [x, y] = xy - yx.

Let J(t) be a Jacobi field along a geodesic δ of Σ . That is, J is a solution of the differential equation

$$\frac{D^2 J}{dt^2} + R_{\delta}(J, \dot{\delta})\dot{\delta} = 0. \tag{7}$$

Next we show that the norm of a Jacobi field is convex. If $x, y \in \mathcal{A}_h$ are regarded as tangent vectors of Σ at the point a, then the following condition (which is a non positive sectional curvature condition) holds:

$$< R_a(x,y)y, x>_a = \tau(R_a(x,y)ya^{-1}xa^{-1}) \le 0.$$

The proof of this fact is straightforward. Then

$$\frac{d^2}{dt^2} \langle J, J \rangle_{\gamma} = 2 \left\{ \langle \frac{D^2 J}{dt^2}, J \rangle_{\gamma} + \langle \frac{DJ}{dt}, \frac{DJ}{dt} \rangle_{\gamma} \right\} =$$

$$= 2 \left\{ -\langle R_{\gamma}(J, \dot{\gamma})\dot{\gamma}, J \rangle_{\gamma} + \langle \frac{DJ}{dt}, \frac{DJ}{dt} \rangle_{\gamma} \right\} \ge 0.$$

In other words, the smooth function $t \mapsto \langle J, J \rangle_{\gamma}$ is convex. We shall need convexity of the *norm* of the Jacobi field (and not of the *square* of the norm just noted).

Proposition 3.3. Let γ be a geodesic of Σ and J a Jacobi field along γ . The real map $t \mapsto \langle J, J \rangle_{\gamma}^{\frac{1}{2}}$ is convex.

Proof. Clearly, is suffices to prove this assertion for a field J which does not vanish. As in Theorem 1 of [3], by the invariance of the connection and the metric under the action of $G_{\mathcal{A}}$, it suffices to consider the case of a geodesic $\gamma(t) = e^{tx}$ starting at $1 \in \Sigma$ $(x \in \mathcal{A}_h)$. For the field $K(t) = e^{-tx/2}J(t)e^{-tx/2}$ the Jacobi equation translates into

$$4\ddot{K} = Kx^2 + x^2K - 2xKx. (8)$$

Moreover

$$_{\gamma}^{\frac{1}{2}} = \tau(\gamma^{-1}J\gamma^{-1}J)^{\frac{1}{2}} = \tau(K^2)^{\frac{1}{2}} = ||K||_2.$$

Let us prove therefore that the map $t \mapsto f(t) = ||K(t)||_2$ is convex, for any (non vanishing) solution K of (8). Note that f(t) is smooth, and $\dot{f} = \tau(K^2)^{-\frac{1}{2}}\tau(K\dot{K})$. Then

$$\ddot{f} = -\tau(K^2)^{-\frac{3}{2}}\tau(K\dot{K})^2 + \tau(K^2)^{-\frac{1}{2}}\{\tau(\dot{K}^2) + \tau(K\ddot{K})\}.$$

Let us multiply this expresion by $\tau(K^2)^{\frac{3}{2}}$ to obtain

$$-\tau (K\dot{K})^{2} + \tau (K^{2})\tau (\dot{K}^{2}) + \tau (K^{2})\tau (K\ddot{K}).$$

The first two terms add up to a non negative number. Indeed, one has $\tau(K\dot{K})^2 \leq \tau(K^2)\tau(\dot{K}^2)$ by the Cauchy-Schwarz inequality for the trace τ . Let us examine the third term $\tau(K^2)\tau(K\ddot{K})$. It suffices to show that $\tau(K\ddot{K})$ is non negative. Using (8),

$$\tau(K\ddot{K}) = \frac{1}{4}\{\tau(K^2x^2) + \tau(Kx^2K) - 2\tau(KxKx)\} = \frac{1}{2}\{\tau(K^2x^2) - \tau(KxKx)\}.$$

This number is positive, again by the Cauchy-Schwarz inequality:

$$\tau(KxKx) = \tau((xK)^*Kx) \le \tau((xK)^*xK)^{\frac{1}{2}}\tau((Kx)^*Kx)^{\frac{1}{2}} = \tau(K^2x^2).$$

Corollary 3.4. If γ and δ are geodesics, the map $f(t) = dist(\gamma(t), \delta(t))$ is a convex function of t.

Proof. As in Theorem 2 of [3], distance between $\gamma(t)$ and $\delta(t)$ is given by the geodesic $\alpha_t(s)$ obtained moving the s variable in a geodesic square h(s,t) with vertices $\gamma(t_0), \delta(t_0), \gamma(t_1), \delta(t_1)$. Taking the partial derivative along the s direction gives a Jacobi field J(s,t) along the geodesic $\beta_s(t) = h(s,t)$ and it also gives the speed of α_t . Hence

$$f(t) = \int_0^1 \|\frac{\partial \alpha_t}{\partial s}(s)\|_{\alpha_t(s)} ds = \int_0^1 \|J(s,t)\|_{h(s,t)} ds.$$

This equation says that f(t) can be written as the limit of a convex combination of convex functions $u_i(t) = ||J(s_i, t)||_{h(s_i, t)}$, so f must be convex itself.

Lemma 3.5. The following inequality holds for any $x, y \in A_h$:

$$dist(e^x, e^y) = \|\ln\left(e^{x/2}e^{-y}e^{x/2}\right)\|_2 \ge \|x - y\|_2.$$
(9)

Proof. Take $\gamma(t) = e^{tx}$, $\delta(t) = e^{ty}$ and f as in the previous corollary. Note that f(0) = 0, hence $f(t)/t \le f(1)$ for any $0 < t \le 1$; hence $\lim_{t \to 0^+} f(t)/t \le f(1)$. Now note that

$$f(t)/t = \frac{1}{t} \|\ln\left(e^{tx/2}e^{-ty}e^{tx/2}\right)\|_{2} = \tau \left(\left[\frac{1}{t}\ln\left(e^{tx/2}e^{-ty}e^{tx/2}\right)\right]^{2}\right)^{\frac{1}{2}}.$$

Since $\lim_{t\to 0^+} \frac{1}{t} \ln\left(e^{tx/2}e^{-ty}e^{tx/2}\right) = \frac{d}{dt} \mid_{t=0} \ln\left(e^{tx/2}e^{-ty}e^{tx/2}\right)$, and the logarithm of $\beta(t) = e^{tx/2}e^{-ty}e^{tx/2}$ can be approximated uniformly by polinomials $p_n(\beta) = \sum_k \alpha_{n,k}\beta^k$ for t close enough to zero (note that $\beta(0) = 1$), and $\frac{d}{dt}\beta\mid_{t=0} = x - y$, we have the desired inequality.

Using the inner product in each tangent space, we can talk about angles between curves and more general subsets of Σ in a natural way; in particular we have:

Lemma 3.6. The sum of the inner angles of any geodesic triangle in the manifold Σ is less or equal than π

Proof. Squaring both sides of inequality (9) leads (by the invariance of the metric for the action of G_A) to

$$l_i^2 \ge l_{i+1}^2 + l_{i-1}^2 - 2l_{i+1}l_{i-1}\cos(\alpha_i)$$

where l_i are the sides of any geodesic triangle and α_i is the angle opposite to l_i . These inequalities say that we can construct an Euclidian comparison triangle in the affine plane with sides l_i ; they also say that the angle β_i (opposite to l_i for this flat triangle) is bigger than α_i . Adding the three angles we have $\alpha_1 + \alpha_2 + \alpha_3 \le \beta_1 + \beta_2 + \beta_3 = \pi$.

4 Convex sets

We are interested in the convex subsets of Σ , that is, subsets $M \subset \Sigma$ such that the ambient geodesic joining two points in M stays in M for any value of t. Note that the simplest of such objects are the geodesics.

It's not hard to see that when two elements $a, b \in \Sigma$ commute, the geodesic triangle spanned by a, b and 1 is convex, hence there is a flat surface containing a, b and 1; indeed, the triangle in Σ is the image of the plane triangle with vertices $0, e_1, e_2$ under the map $T : \mathbb{R}^2 \to \Sigma$ given by

$$T(x,y) = e^{x\ln(a) - y\ln(b)}$$

In particular the geodesic joining a and b is the image of the segment

$$\{(x,y): x,y \ge 0, x+y=1\}.$$

This is not true in the general case (though the length of the segment is a lower bound for the length of that geodesic, as Lemma 9 shows).

Definition 4.1. An exponential set $M \subset \Sigma$ is the exponential of a (closed, self-adjoint) subspace through the origin. In other words, $M = e^H$ with H a closed subspace of A_h .

Lemma 4.2. If M is a convex exponential set in Σ , the geodesic symmetry σ_p : $q \mapsto pq^{-1}p$ maps M into M for any $p \in M$

Proof. The map σ_p maps any geodesic through $p, \gamma(t) = p^{\frac{1}{2}} e^{tp^{-\frac{1}{2}} vp^{-\frac{1}{2}}} p^{\frac{1}{2}}$ onto $\gamma(-t)$; now it is clear that it is an isommetry of Σ and it maps M into M.

Note that, if M is convex and $a \in M$, then $a^{\alpha} = e^{\alpha \ln a}$ is in M for any real α . This observation together with the previous lemma leads to the following characterization of convexity:

Proposition 4.3. If M is a convex, exponential set in Σ , then

$$aba \in M \text{ whenever } a, b \in M$$
 (10)

Proof. Note that
$$aba = a^{\frac{3}{2}} \left(a^{\frac{1}{2}} b^{-1} a^{\frac{1}{2}} \right)^{-1} a^{\frac{3}{2}} = \sigma_{a^{\frac{3}{2}}} \circ \sigma_{a^{\frac{1}{2}}}(b).$$

The converse of this last statement is also true (this can be easily seen iterating the property above in order to construct $\gamma(t)$ for given $p, q \in M$).

Let us see how this property looks in the tangent H (recall that $M = e^H$). This result is related to the results of [7] by H. Porta and L. Recht:

Theorem 4.4. If H is a closed subspace of A_h (in the norm topology of A), then $M = e^H$ is a geodesically convex subset of Σ if and only if $[x, [x, y]] \in H$ for any $x, y \in H$.

Proof. We use property (10) above to identify convex sets; the proof follows the guidelines of [5] for matrices, and we translate it here.

We first assume H has the double bracket property. Set $D_x : A \to A$, $D_x = L_x - R_x$, the difference between left and right multiplication by x in A.

Let's consider the completion of \mathcal{A} with respect to the trace, namely $\mathcal{H} = L^2(\mathcal{A}, \tau)$. Clearly $\mathcal{H}_{\mathbb{R}} = L^2_{\mathbb{R}}(\mathcal{A}_h, \tau)$ contains as a proper, closed subspace the completion of H, namely $\mathcal{H}_1 = L^2_{\mathbb{R}}(H, \tau)$. Since τ is a normal trace, the involution * extends to a bounded antilinear operator J of $\mathcal{H}_{\mathbb{R}}$, and the map D_x extends uniquely to a bounded linear operator of $\mathcal{H}_{\mathbb{R}}$ (which we will still call D_x).

First we establish the identity

$$T_x(y) = g(D_x/2)(y) \tag{11}$$

where T_x is the extension of the map from Corollary 2.2, and

$$g(z) := \frac{\sinh(z)}{z} = \sum_{n>0} \frac{z^{2n}}{(2n+1)!}$$

is an entire function. Note that $g(z) = (2z)^{-1}(e^z - e^{-z})$. To prove (11), we take derivative with respect to t in the identity $X(t)e^{X(t)} = e^{X(t)}X(t)$, where X(t) = x + ty; after rearranging the terms we come up with

$$\frac{1}{2} \left(e^{D_x/2} - e^{-D_x/2} \right) y = (D_x \circ T_x)(y).$$

Note that if D_x were invertible, we would be set; this is not necessarily the case. However, $D_x^2 = D_x \circ D_x$ is selfadjoint when restricted to \mathcal{A} , and since T_x (more precisely, its extension) is also sefaldjoint (Corollary 2.2), the operator $T = T_x - g(D_x/2)$ is selfadjoint on \mathcal{A} , and hence on \mathcal{H} (note that g(z) involves only even powers of z). The equation above says that we have proved that $(D_x \circ T)(y) = 0$ for any $y \in \mathcal{A}$; in other words T maps \mathcal{H} into $\{x\}' = \{b \in \mathcal{A}_h : bx = xb\}$. A straightforward computation shows that Tb = 0 for any $b \in \{x\}'$, which proves that T = 0, i.e. equation (11) holds.

Now, for $x, y \in H$ consider the curve $e^{\alpha(t)} = e^{tx}e^{y}e^{tx}$. Clearly $\alpha(0) = y \in H \subset \mathcal{H}_1$; we will prove that α obeys a differential equation in $\mathcal{H}_{\mathbb{R}}$ which has a flow that maps \mathcal{H}_1 into \mathcal{H}_1 , and by the uniqueness of the solution of such equation we will have $e^x e^y e^x = e^{\alpha(1)} \in e^H = M$.

Differentiating at $t = t_0$ the equation yields to

$$xe^{\alpha(t_0)} + e^{\alpha(t_0)}x = dexp_{\alpha(t_0)}(\dot{\alpha}(t_0)) = e^{\alpha(t_0)/2}T_{\alpha(t_0)}(\dot{\alpha}(t_0))e^{\alpha(t_0)/2} =$$
$$= e^{\alpha(t_0)/2}g(D_{\alpha(t_0)}/2)(\dot{\alpha}(t_0))e^{\alpha(t_0)/2}.$$

Note that g(z) is invertible whenever z is a bounded linear operator, and also that the power series for z coth(z/2) involves only even powers of z. On the other hand, $D_z/2 = D_{z/2}$ and $e^z x e^{-z} = e^{D_z} x$, hence

$$\dot{\alpha} = g^{-1}(D_{\alpha}/2) \circ (e^{-\alpha/2}xe^{\alpha/2} + e^{\alpha/2}xe^{-\alpha/2}) =$$

$$= g^{-1}(D_{\alpha/2}) \circ (e^{D_{\alpha}/2} + e^{-D_{\alpha}/2})(x) = D_{\alpha} \coth(D_{\alpha/2})(x) = \sum_{n} c_{n} D_{\alpha}^{2n} x =$$

$$= \sum_{n} c_{n} D_{\alpha}^{2} \circ \cdots \circ D_{\alpha}^{2}(x) = F(\alpha).$$

Since $D_z^2(x) = [z, [z, x]]$, $F(z) = \sum_n c_n D_z^{2n}(x)$ can be regarded as a map from \mathcal{H}_1 to \mathcal{H}_1 , and since it is clearly an analytic map of \mathcal{H} into \mathcal{H} , it fulfills a Lipschitz

condition. Now the unique solution must be $\alpha(t) = \ln(e^{tx}e^{y}e^{tx})$. This proves that M is convex whenever H has the double bracket property.

To prove the other implication, assume $M = e^H$ is convex and H is closed in the norm topology of \mathcal{A}_h . Clearly the path α stays in H for any value of t (here $e^{\alpha(t)} = e^{tx}e^ye^{tx}$), and the same is true for $\dot{\alpha}$. Now since $\dot{\alpha}(t) = D_{\alpha(t)} \coth(D_{\alpha(t)}/2)x$, we have

$$\lim_{t \to 0} \frac{\dot{\alpha}(t) - \dot{\alpha}(0)}{t^2} = \lim_{t \to 0} \frac{\left(1 + \frac{1}{12}t^2D_{\alpha(t)}^2\right)x - x}{t^2} + tO(t) = \frac{1}{12}D_y^2(x)$$

which proves that $D_y^2(x) = [y, [y, x]]$ belongs to H whenever x and y are in H. \square

5 Projections

From now on assume $M=e^H$ is a convex exponential set in Σ . As before, we identify the derivatives of all the geodesics at $p \in M$ with the tangent space of M at p, in order to define the angles between curves and sets in a natural way: note that in this way, $T_1M=H$ and $T_pM=p^{\frac{1}{2}}Hp^{\frac{1}{2}}$ (which can be thought of as the parallel transport along the geodesic joining 1 and p in M).

In particular, $T_1\Sigma$ is naturally identified with \mathcal{A}_h and the same is true for $T_p\Sigma$, for any $p \in \Sigma$, since $p^{\frac{1}{2}}\mathcal{A}_h p^{\frac{1}{2}} = \mathcal{A}_h$ (this is clear also from the fact that Σ is open in \mathcal{A}_h).

Lemma 5.1. Let $r \in \Sigma$. There is at most one point $p = \Pi_M(r)$ in M such that the geodesic joining r and p is orthogonal to M at p

Proof. Assume there are two points p and p' in M and two vectors v and v' orthogonal to M at p and p' respectively such that $\gamma_1(1) = Exp_p(v) = \gamma_2(1) = Exp_{p'}(v') := r$, and consider the geodesic triangle with sides the given geodesics and the unique geodesic in M joining p and p'. Since the angles at p and p' are right angles, and the summ of the inner angles of any such geodesic triangle is less or equal than π , it must be that the angle at r is zero: since geodesics are unique (given an initial velocity and an initial position r), it must be that $\gamma_1 = \gamma_2$, hence p = p' and v = v'.

Set NM as the normal bundle of M, i.e. $NM = \{(p, v) : p \in M, v \in (T_pM)^{\perp}\}$. Consider the map $E : NM \to \Sigma$ given by $(p, v) \mapsto Exp_p(v)$; since E is analytic and with the right identifications has differential (at (p, 0)) the identity map, E(NM) contains an open neighbourhood of M in Σ (with the norm topology).

Lemma 5.2. The map $\Pi_M : E(NM) \to M$ that assigns the endpoint of the minimizing geodesic is contractive for the geodesic metric.

Proof. If r, s are two points in E(NM), we will prove that this projection is contractive. Assume $\Pi_M(r) = p, \Pi_M(s) = q \in M, v \in T_p\Sigma$ is orthogonal to M at p and w is orthogonal to M at q; let's consider the distance function

$$f(t) = dist^{2}(Exp_{p}(tv), Exp_{q}(tw)) = dist^{2}(\gamma_{1}(t), \gamma_{2}(t))$$

where $\gamma_1(t)$ is the only geodesic with initial velocity v starting at p and γ_2 is the only geodesic with initial speed w starting at q. Namely,

$$\gamma_1(t) = p^{\frac{1}{2}} e^{tp^{-\frac{1}{2}} vp^{-\frac{1}{2}}} p^{\frac{1}{2}}$$
 and $\gamma_2(t) = q^{\frac{1}{2}} e^{tq^{-\frac{1}{2}} wq^{-\frac{1}{2}}} q^{\frac{1}{2}}.$

Since $v \in (T_n M)^{\perp}$ and $w \in (T_n M)^{\perp}$, we have

$$\langle v, p^{\frac{1}{2}} x p^{\frac{1}{2}} \rangle_{p} = \tau \left(x p^{-\frac{1}{2}} v p^{-\frac{1}{2}} \right) = 0 \text{ for any } x \in H = T_{1}M \text{ and}$$

$$\langle w, q^{\frac{1}{2}} y q^{\frac{1}{2}} \rangle_{q} = \tau \left(y q^{-\frac{1}{2}} w q^{-\frac{1}{2}} \right) = 0 \text{ for any } y \in H = T_{1}M.$$
(12)

Now we use the formula $dist(e^A, e^B) = \|\ln(e^{A/2}e^{-B}e^{A/2})\|_2$ for $A = \ln(\gamma_1(t))$ and $B = \ln(\gamma_2(t))$, to write

$$f(t) = \|\ln(\gamma_1^{\frac{1}{2}}\gamma_2^{-1}\gamma_1^{\frac{1}{2}})\|_2^2 = \tau \left(\ln^2(\gamma_1^{\frac{1}{2}}\gamma_2^{-1}\gamma_1^{\frac{1}{2}})\right).$$

Assume that C is a simple, positively oriented curve in \mathbb{C} , around the spectrum of $\alpha_0 = p^{\frac{1}{2}}q^{-1}p^{\frac{1}{2}}$. Then we can use the Cauchy formula to calculate $\ln^2(a)$ for any element $a \in \mathcal{A}$ such that $\sigma(a) \subset int(C)$, namely

$$\ln^2(a) = \frac{1}{2\pi i} \int_C \ln^2(z) (z - a)^{-1} dz. \tag{13}$$

Naming $\alpha(t) = \gamma_1^{\frac{1}{2}}(t)\gamma_2^{-1}(t)\gamma_1^{\frac{1}{2}}(t)$, this formula holds true for $\alpha_0 = \alpha(0)$ and for $\alpha(t)$ for t sufficiently small. Note that

$$f(t) = \tau \left(\gamma^{-\frac{1}{2}}(t) \gamma^{\frac{1}{2}}(t) \ln^2(\alpha(t)) \right) = \tau \left(\gamma^{-\frac{1}{2}}(t) \ln^2(\alpha(t)) \gamma^{\frac{1}{2}}(t) \right).$$

If x is invertible in \mathcal{A} , $xg(a)x^{-1} = g(xax^{-1})$ for any element $a \in \mathcal{A}$ and any analytic function g in a neighbourhood of $\sigma(a)$. Then

$$f(t) = \tau \left(\ln^2 \left[\gamma_1(t) \gamma_2^{-1}(t) \right] \right) = \frac{1}{2\pi i} \int_C \ln^2(z) \, \tau \left[\left(z - \gamma_1(t) \gamma_2^{-1}(t) \right)^{-1} \right] \, dz.$$

Now we compute f'(0); note first that $\gamma_1(0)\gamma_2^{-1}(0) = pq^{-1}$ and also that

$$\frac{d}{dt} {}_{t=0} \gamma_1(t) \gamma_2^{-1}(t) = -vq^{-1} + pq^{-1}wq^{-1}.$$

Using the properties of the trace we get

$$\frac{d}{dt}_{t=0}f(t) = -\frac{1}{2\pi i} \int_C \ln^2(z) \, \tau \left[\left(z - pq^{-1} \right)^{-2} \left(-vq^{-1} + pq^{-1}wq^{-1} \right) \right] \, dz =
= \tau \left[\left(-\frac{1}{2\pi i} \int_C \ln^2(z) \, \left(z - pq^{-1} \right)^{-2} \, dz \right) \left(-vq^{-1} + pq^{-1}wq^{-1} \right) \right].$$

If we integrate by parts the first factor inside the trace, we obtain (note that $\frac{d}{dz}ln^2(z) = 2\ln(z)z^{-1} = 2z^{-1}\ln(z)$ and C is a closed curve) that

$$\begin{split} \dot{f}(0) &= \tau \left[\left(\frac{1}{2\pi i} \int_C 2 \ln(z) z^{-1} \, \left(z - p q^{-1} \right)^{-1} dz \right) \left(-v q^{-1} + p q^{-1} w q^{-1} \right) \right] = \\ &= -\tau \left[\left(\frac{1}{2\pi i} \int_C 2 \ln(z) z^{-1} \, \left(z - p q^{-1} \right)^{-1} dz \right) v q^{-1} \right] + \\ &+ \tau \left[\left(\frac{1}{2\pi i} \int_C 2 \ln(z) z^{-1} \, \left(z - p q^{-1} \right)^{-1} dz \right) p q^{-1} w q^{-1} \right]. \end{split}$$

Therefore,

$$\dot{f}(0) = -\frac{1}{2\pi i} \int_{C} 2\ln(z) z^{-1} \tau \left[q^{-1} \left(z - pq^{-1} \right)^{-1} v \right] dz +$$

$$+ \frac{1}{2\pi i} \int_{C} 2\ln(z) z^{-1} \tau \left[q^{-\frac{1}{2}} \left(z - pq^{-1} \right)^{-1} pq^{-1} wq^{-\frac{1}{2}} \right] dz.$$

Using the elementary identities

$$p^{\frac{1}{2}}(z-p^{\frac{1}{2}}q^{-1}p^{\frac{1}{2}})^{-1}p^{-\frac{1}{2}}=(z-pq^{-1})^{-1}=q^{\frac{1}{2}}(z-q^{-\frac{1}{2}}pq^{-\frac{1}{2}})^{-1}q^{-\frac{1}{2}}$$

one arrives to the expression

$$\begin{split} \dot{f}(0) &= -\frac{1}{2\pi i} \int_C 2 \ln(z) z^{-1} \, \tau \left[q^{-1} p^{\frac{1}{2}} \left(z - p^{\frac{1}{2}} q^{-1} p^{\frac{1}{2}} \right)^{-1} p^{-\frac{1}{2}} v \right] dz + \\ &\quad + \frac{1}{2\pi i} \int_C 2 \ln(z) z^{-1} \, \tau \left[\left(z - q^{-\frac{1}{2}} p q^{-\frac{1}{2}} \right)^{-1} q^{-\frac{1}{2}} p q^{-\frac{1}{2}} q^{-\frac{1}{2}} w q^{-\frac{1}{2}} \right] dz = \\ &\quad = -2\tau \left[q^{-1} p^{\frac{1}{2}} p^{-\frac{1}{2}} q p^{-\frac{1}{2}} \ln(p^{\frac{1}{2}} q^{-1} p^{\frac{1}{2}}) p^{-\frac{1}{2}} v \right] + \\ &\quad + 2\tau \left[\ln(q^{-\frac{1}{2}} p q^{-\frac{1}{2}}) q^{\frac{1}{2}} p^{-1} q^{\frac{1}{2}} q^{-\frac{1}{2}} p q^{-\frac{1}{2}} q^{-\frac{1}{2}} w q^{-\frac{1}{2}} \right) = \\ &\quad = -2\tau \left[\ln(p^{\frac{1}{2}} q^{-1} p^{\frac{1}{2}}) p^{-\frac{1}{2}} v p^{-\frac{1}{2}} \right] + 2\tau \left[\ln(q^{-\frac{1}{2}} p q^{-\frac{1}{2}}) q^{-\frac{1}{2}} w q^{-\frac{1}{2}} \right] = 0 + 0 = 0, \end{split}$$

which holds by the orthogonality relations (12), naming $x = \ln(p^{\frac{1}{2}}q^{-1}p^{\frac{1}{2}})$ (recall that M is convex), and $y = \ln(q^{-\frac{1}{2}}pq^{-\frac{1}{2}})$.

Since f(t) is a convex function, f has a global minimum at t = 0, which proves that $dist(\Pi_M(r), \Pi_M(s)) = dist(p, q) \leq dist(r, s)$.

Lemma 5.3. Let $r \in \Sigma$ and let M be a convex exponential set. The there exists $p \in M$ such that dist(r, M) = dist(r, p) if and only if there is a geodesic through r orthogonal to M.

Proof. Assume first that there is a point $p \in M$ such that the geodesic γ through p and r is orthogonal to M at p. Now for any point $q \in M$, take a geodesic β joining q and r, and a geodesic δ joining p to q. Consider the geodesic triangle with sides γ, β, δ ; the angle opposite to β is a right angle, so (see Lemma 3.6)

$$Length(\beta)^2 \ge Length(\delta)^2 + Length(\gamma)^2 \ge Length(\gamma)^2$$
.

This proves that $dist(p,r) \leq dist(q,r)$ for any $q \in M$.

Assume now that $p \in M$ has the minimizing property and consider, for any point $q \in M$, the geodesic $\gamma_{p,q}(s)$ joining p to q (note that it is inside M for any s by virtue of the convexity). Now consider the family of geodesics

$$\gamma_s(t) = \gamma_{r,\gamma_{pq}(s)}(t) = r^{\frac{1}{2}} \left(r^{-\frac{1}{2}} p^{\frac{1}{2}} \left(p^{-\frac{1}{2}} q p^{-\frac{1}{2}} \right)^s p^{\frac{1}{2}} r^{-\frac{1}{2}} \right)^t r^{\frac{1}{2}}$$

that is, the family of geodesics joining r to $\gamma_{p,q}(s)$.

Put $g(s) = Length(\gamma_s)^2 = dist(r, \gamma_{p,q}(s))$. This function has a minimum at s = 0, hence (since it is C^{∞}) it must be that $\dot{g}(0) = 0$. As in the proof of the previous theorem, we have

$$g(s) = \tau \left(\ln^2(r\gamma_{p,q}(s)^{-1}) \right) = \frac{1}{2\pi i} \int_C \ln^2(z) \, \tau \left[\left(z - r\gamma_{p,q}^{-1}(s) \right)^{-1} \right] \, dz.$$

Taking the derivative at s=0 and integrating by parts we obtain

$$0 = \dot{g}(0) = -2\tau \left(\ln(rp^{-1})\ln(qp^{-1})p^{-1}\right) = 2\tau \left(\ln(qp^{-1})p^{-1}\ln(pr^{-1})\right).$$

On the other hand, the angle subtended by $\gamma_{p,q}$ and $\gamma_{r,p}$ at p is

$$<\dot{\gamma}_{r,p}(1),\dot{\gamma}_{p,q}(0)>_p=\tau\left(\ln(qp^{-1})p^{-1}\ln(pr^{-1})\right).$$

This proves that $\gamma_{r,p}$ is orthogonal to any geodesic at p contained in M, and by definition, it is orthogonal to M.

The following is related to the main result in [6] by H. Porta and L. Recht:

Theorem 5.4. If $M = e^H$ is a convex exponential set, and there is a closed, orthogonal supplement S for H (namely $A_h = H \oplus_{\perp_{\tau}} S$) then for any point $r \in \Sigma$ there is a geodesic through r orthogonal to M.

Proof. Exactly as in [6], there is an equality of sets $E(NM) = \Sigma$, where NM stands for the normal bundle of M, i.e the pairs (p, v) with $p \in M$ and $v \perp_p M$.

The typical examples for this situation arise when $H = B_h$ for a subalgebra B of A. In this case, by a result of Takesaki [8], there is a conditional expectation $\mathcal{E} : A \to A$ with rank B, compatible with τ (i.e $\tau(\mathcal{E}(x)) = \tau(x)$ for any $x \in A$).

Corollary 5.5. If H is a closed subspace in A_h (supplemented as in the previous theorem) such that $[x, [x, y]] \in H$ whenever $x, y \in H$, then for any $z \in A_h$ we can factor

$$e^z = e^y e^w e^y$$

for unique $y \in H$ and $w \in A_h$ such that $\tau(wx) = 0$ for any $x \in H$. Moreover, e^{2y} minimizes (geodesic) distance between $M = e^H$ and e^z , and is unique with that property.

Corollary 5.6. Fix \mathcal{D} a m.a.s.a of \mathcal{A} . Then for any $x \in \mathcal{A}_h$ there are unique $d \in \mathcal{D}_h$ and $v \in \mathcal{A}_h$ such that $\tau(vz) = 0$ for any $z \in \mathcal{D}$, and $e^x = de^v d$

Corollary 5.7. If H is a closed, supplemented subspace in A_h such that $[x, [x, y]] \in H$ whenever $x, y \in H$, then for any $g \in G_A$ we can factor

$$g = e^x e^y u$$

for unique $x \in H, y \in H^{\perp}$ and u in the unitary group of A.

Proof. Note that $gg^* \in \Sigma$ hence $gg^* = e^x e^{2y} e^x$ where x, y are as required. Now take $u = e^{-y} e^{-x} g$; a straightforward computation shows that $uu^* = u^* u = 1$. Uniqueness follows from the uniqueness of x, y.

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