Projections with fixed difference: a Hopf-Rinow theorem.

Esteban Andruchow, Gustavo Corach, Lázaro Recht

May 17, 2018

Abstract

The set $\mathcal{D}_{A_0}$, of pairs of orthogonal projections (P, Q) in generic position with fixed difference $P - Q = A_0$, is shown to be a homogeneus smooth manifold: it is the quotient of the unitary group of the commutant $\{A_0\}'$ divided by the unitary subgroup of the commutant $\{P_0, Q_0\}'$, where (P_0, Q_0) is any fixed pair in $\mathcal{D}_{A_0}$. Endowed with a natural reductive structure (a linear connection) and the quotient Finsler metric of the operator norm, it behaves as a classic Riemannian space: any two pairs in $\mathcal{D}_{A_0}$ are joined by a geodesic of minimal length. Given a base pair (P_0, Q_0), pairs in an open dense subset of $\mathcal{D}_{A_0}$ can be joined to (P_0, Q_0) by a *unique* minimal geodesic.

2010 MSC: 58B20, 47B15.
Keywords: Projections, pairs of projections, differences of projections.

1 Introduction

Let $\mathcal{H}$ be a Hilbert space, denote by $\mathcal{B}(\mathcal{H})$ the algebra of bounded linear operators in $\mathcal{H}$, and by $\mathcal{P}(\mathcal{H})$ the set of (orthogonal) projections in $\mathcal{H}$. We study here the class $\mathcal{D}$ of operators which are differences of projections,

$$\mathcal{D} = \{P - Q : P, Q \in \mathcal{P}(\mathcal{H})\},$$

and for $A \in \mathcal{D}$, the set

$$\mathcal{D}_A = \{(P, Q) \in \mathcal{P}(\mathcal{H}) \times \mathcal{P}(\mathcal{H}) : P - Q = A\}.$$

Results on differences $P - Q$ appeared since the 1940's, as part of the *two subspaces problem*: to find a complete set of unitary invariants for a pair of closed subspaces M, N (or equivalently, for a pair of projections P, Q). This problem was solved by J. Dixmier [15], who obtained a characterization of $\mathcal{D}$. An operator $A \in \mathcal{B}(\mathcal{H})$ belongs to $\mathcal{D}$ if and only if $A^* = A$, $\|A\| \leq 1$ and there exists a symmetry V in $\mathcal{H}' = N(A^2 - 1)^{\perp}$ such that $AV = VA$ in $\mathcal{H}'$ (a *symmetry* is a selfadjoint unitary operator). This form of Dixmier's result is due to Ch. Davis [13], who found a nice solution of the two subspaces problem by a geometric study of the closeness and separation operators of a pair P, Q: $C(P, Q) = PQP + (1 - P)(1 - Q)(1 - P)$ is called the closeness operator of P, Q, and $S(P, Q) = P(1 - Q)P + (1 - P)Q(1 - P)$ is the separation operator of P, Q. Observe that, if $A = P - Q$, then $C = 1 - A^2$ and $S = A^2$. In [9] J. Avron, R. Seiler and B. Simon defined and studied Fredholm pairs of projections, and an index for them: (P, Q) is a Fredholm

pair if $P|_{R(Q)} : R(Q) \to R(P)$ is a Fredholm operator, whose index is called the index of the pair. Their methods rely on an extensive use of the differences $A = P - Q$ and $B = P + Q - 1$. For a nice presentation of these results, see W. Amrein and K. Sinha [1]

A more recent study of $\mathcal{D}$ can be found in [2], where several known facts on the differential geometry of $\mathcal{P}(\mathcal{H})$ were used to describe, for instance, the interior and boundary of $\mathcal{D}$, its connected components, and also some special parts of $\mathcal{D}$ (elements in $\mathcal{D}$ which are Fredholm, compact, or nuclear).

In [30], W. Shi., G. Ji and H. Du studied several properties of $\mathcal{D}_A$, for any $A \in \mathcal{D}$. In particular, they proved that $\mathcal{D}_{A_0} \subset \mathcal{B}(\mathcal{H}_0)$ is connected, where $\mathcal{H}_0 = \{N(A^2 - 1) \oplus N(A)\}^{\perp}$ and $A_0 = A|_{\mathcal{H}_0}$.

The main goal of this paper is to present $\mathcal{D}_{A_0}$ as a homogeneous space and a differentiable manifold. As such, following ideas of Durán, Mata-Lorenzo and Recht [17], the space $\mathcal{D}_{A_0}$ has a natural invariant Finsler metric. Also, using a well known characterization by Halmos [21], of pairs of projections in generic position, we show that $\mathcal{D}_{A_0}$ has a reductive structure, a fact which enables one to introduce a linear connection in this space, and to compute its geodesics (given by one-parameter unitary groups acting on a given pair (P_0, Q_0)). We show that with the Finsler metric and the reductive structure, $\mathcal{D}_{A_0}$ satisfies a Hopf-Rinow theorem: pairs in $\mathcal{D}_{A_0}$ are joined by a geodesic of minimal length. Moreover, on a dense open subset of $\mathcal{D}_{A_0}$, such geodesic is unique.

In Section 2, we present Davis' characterization of $\mathcal{D}$ by means of the *Halmos decomposition* of $\mathcal{H}$ (in the presence of a pair $P, Q \in \mathcal{P}(\mathcal{H})$). Using Davis' and Halmos' tools, we show that the Friedrich's angle is constant in $\mathcal{D}_A$. Recall (see Deutsch [14]) that $\alpha_F(M, N) \in [0, \pi/2]$ is the Friedrich's angle between the closed subspaces M, N if

$$\cos(\alpha_F(M, N)) = \sup\{|\langle \mu, \nu \rangle| : \mu \in M \ominus N, \nu \in N \ominus M, \|\mu\| = \|\nu\| = 1\} = \|P_M P_N - P_{M \cap N}\|.$$

Moreover, it is shown that $\cos(\alpha_F(M, N)) = \|P_0 Q_0\|$ (=constant) for any $P = P_M$, $Q = P_N$ such that $P - Q = A$, where P_0, Q_0 denote the reductions of P, Q to the common invariant subspace $\mathcal{H}_0 = \{N(A) \oplus N(A^2 - 1)\}^{\perp}$. Hereafter, P_0, Q_0, A_0 will be called the *generic part* of P, Q, A, respectively. Also in this section we show that, with the usual order of positive definite operators, the set $\{P_0 + Q_0 : (P_0, Q_0) \in \mathcal{D}_{A_0}\}$ cannot be ordered: $P_0 + Q_p \le P_0' + Q_0'$ if and only if $P_0 = P_0'$ and $Q_0 = Q_0'$. In Section 3 we introduce the action of the unitary group $\mathcal{U}_{\mathcal{A}}$ of

$$\mathcal{A} = \{A\}' = \{T \in \mathcal{B}(\mathcal{H}) : TA = AT\}$$

on $\mathcal{D}_A$. $\mathcal{A}$ and $\mathcal{U}_{\mathcal{A}}$ are also reduced by $\mathcal{H}_0$. It is proven that the generic part of $\mathcal{U}_{\mathcal{A}}$ acts transitively on $\mathcal{D}_{A_0}$, and from this follows that $\mathcal{D}_{A_0}$ is connected (as proved by Shi, Ji and Du in [30]). Section 4 contains a description of $\mathcal{A}_0$ in terms of Halmos' decomposition, which will be used later. In Section 5 we present some examples in $\mathcal{D}$. In Section 6 we show that $\mathcal{D}_{A_0}$ is a differentiable homogeneous manifold, with a natural reductive structure. For instance, the geodesic curves can be computed, and we show that the exponential map of the linear connection is surjective. We endow the tangent spaces of $\mathcal{D}_{A_0}$ with the quotient norm, as defined by Durán, Mata-Lorenzo and Recht in [17], and show that with this metric, the geodesics of the reducive connection are minimal up to the border of $\mathcal{D}_{A_0}$.

2 Davis' characterization

If $T \in \mathcal{B}(\mathcal{H})$, denote by $R(T)$ and $N(T)$ the range and nullspace of T, respectively. If $A \in \mathcal{D}$, the space $\mathcal{H}$ can be decomposed orthogonally as

$$\mathcal{H} = N(A) \oplus N(A^2 - 1) \oplus \mathcal{H}_0, \tag{1}$$

where $\mathcal{H}_0 = \left(N(A) \oplus N(A^2 - 1)\right)^{\perp}$. Note that $N(A^2 - 1) = N(A - 1) \oplus N(A + 1)$. For any presentation $A = P - Q$, it is straightforward to verify that

$$N(A) = R(P) \cap R(Q) \oplus N(P) \cap N(Q), \ \ N(A-1) = R(P) \cap N(Q) \text{ and } N(A+1) = N(P) \cap R(Q).$$

So that the decomposition (1) is essentially the decomposition considered by Dixmier [15] and Halmos [21] to study the equivalence of pair of projections. In particular, the subspace $\mathcal{H}_0$ is usually called the *generic part* of P and Q, or more properly, the generic part of $A = P - Q$. Therefore, the decomposition (1) reduces simultaneously any pair P, Q in $\mathcal{D}_A$.

Using the decomposition (1), the set $\mathcal{D}_A$ is factorized as follows:

1. In the subspace $N(A^2 - 1) = N(A-1) \oplus N(A+1)$, A is given by $A = 1_{N(A-1)} \oplus -1_{N(A+1)}$. That is, any pair $(P, Q) \in \mathcal{D}_A$ coincides with $(P_{N(A-1)}, P_{N(A+1)})$ in this subspace.

2. In the subspace $N(A)$, the pairs $(P, Q) \in \mathcal{D}_A$ reduce to pairs of the form (P', P'), with $P' \in \mathcal{P}(N(A))$. Thus, if $N(A)$ is non trivial, the structure of $\mathcal{D}_A|_{N(A)}$ is that of $\mathcal{P}(N(A))$.

3. The structure of $\mathcal{D}_A$ in $\mathcal{H}_0$ was characterized by Davis [13]. In Theorem 2.2 below we describe the results obtained by Davis [13] on this set.

A *symmetry* $V \in \mathcal{B}(\mathcal{H})$ is a selfadjoint unitary operator: $V^* = V^{-1} = V$. Symmetries are special cases of difference of projections: $V = P_{+1} - P_{-1}$, where $P_{\pm 1}$ are the orthogonal projections onto the eigenspaces $\{\xi \in \mathcal{H} : V\xi = \pm\xi\}$. Also note that $V = 2P_{+1} - 1$ and $P_{\pm 1} = \frac{1}{2}(1 \pm V)$.

Let us summarize the information above:

Remark 2.1. In the decomposition $\mathcal{H} = N(A) \oplus N(A^2 - 1) \oplus \mathcal{H}_0$, the set $\mathcal{D}_A$ is decomposed as

$$\mathcal{D}_A = \mathcal{P}_{N(A)} \oplus \{A_{\pm 1}\} \oplus \mathcal{D}_{A_0},$$

where $A_{\pm 1} = P_{N(A-1)} - P_{N(A+1)}$ is a symmetry. It follows that $\mathcal{D}_A$ consists of a single element if and only if A is a symmetry.

As announced, let us describe the structure of $\mathcal{D}_{A_0}$:

Theorem 2.2. (essentially [13]) *Let $A \in \mathcal{D}$, and let A_0 be its generic part. There exist one to one correspondences between*

- *Pairs (P_0, Q_0) such that $P_0 - Q_0 = A_0$.*

- *Symmetries V in $\mathcal{H}_0$ such that $VA_0 = -A_0 V$.*

- *Closed subspaces $\mathcal{S}$ of $\mathcal{H}_0$ such that $A_0(\mathcal{S}) \subset \mathcal{S}^{\perp}$ and $A_0(\mathcal{S}^{\perp}) \subset \mathcal{S}$.*

- *Projections $E \in \mathcal{P}(\mathcal{H}_0)$ such that $EA_0E = (1 - E)A_0(1 - E) = 0$.*

Proof. Given a pair $(P_0, Q_0) \in \mathcal{D}_{A_0}$, one obtains a symmetry which anti-commutes with A_0 as follows. Consider the selfadjoint operator $S = P_0 + Q_0 - 1$. Note that $S = P_0 - (1 - Q_0)$ is also a difference of projections. Its nullspace is trivial:

$$N(S) = R(P_0) \cap R(1 - Q_0) \oplus N(P_0) \cap N(1 - Q_0) = R(P_0) \cap N(Q_0) \oplus N(P_0) \cap R(Q_0) = \{0\}.$$

Therefore, the polar decomposition of S, $S = V|S| = |S|V$ yields a symmetry V (note that V is the sign function of S). Clearly, $SP_0 = P_0 Q_0 = Q_0 S$ and $SQ_0 = Q_0 P_0 = P_0 S$. In particular, this implies that S^2 commutes with P_0 and Q_0. Then $|S| = (S^2)^{1/2}$ also commutes with both projections. It follows that $VP_0 = Q_0 V$ and $VQ_0 = P_0 V$. Then

$$VA_0 = VP_0 - VQ_0 = Q_0 V - P_0 V = -A_0 V.$$

Given a symmetry V which anti-commutes with A_0, put (see [13], p. **181**)

$$P_V = \frac{1}{2}\{1 + A_0 + (1 - A_0^2)^{1/2}V\} \quad \text{and} \quad Q_V = \frac{1}{2}\{1 - A_0 + (1 - A_0^2)^{1/2}V\}.$$

Straightforward computations show that $P_V, Q_V \in \mathcal{P}(\mathcal{H}_0)$, $P_V - Q_V = A_0$, and $P_V + Q_V - 1 = (1 - A_0^2)^{1/2}V$. Then, since V and A_0^2 commute,

$$(P_V + Q_V - 1)^2 = 1 - A_0^2, \text{ i.e., } |P_V + Q_V - 1| = (1 - A_0^2)^{1/2},$$

and $P_V + Q_V - 1 = |P_V + Q_V - 1|V$. That is, the correspondence betwen pairs and symmetries is reciprocal.

Given a symmetry V which anti-commutes with A_0, let $\mathcal{S} = \{\xi \in \mathcal{H}_0 : V\xi = \xi\}$, so that $\mathcal{S}^\perp = \{\xi \in \mathcal{H}_0 : V\xi = -\xi\}$. If $\xi \in \mathcal{S}$, $VA_0\xi = -A_0 V\xi = -A_0\xi$, i.e., $A\xi \in \mathcal{S}^\perp$. Similarly, $A(\mathcal{S}^\perp) \subset \mathcal{S}$. The converse holds: if $A_0(\mathcal{S}) \subset \mathcal{S}^\perp$ and $A(\mathcal{S}^\perp) \subset \mathcal{S}$, then the symmetry $V = P_{\mathcal{S}} - P\mathcal{S}^\perp = 2P_{\mathcal{S}} - 1$ anti-commutes with A_0. In fact, if $\xi \in \mathcal{S}$,

$$(2P_{\mathcal{S}} - 1)A_0\xi = 2P_{\mathcal{S}}A_0\xi - A_0\xi = -A_0\xi = -A_0(2P_{\mathcal{S}} - 1)\xi;$$

if $\eta \in \mathcal{S}^\perp$,

$$(2P_{\mathcal{S}} - 1)A_0\eta = 2P_{\mathcal{S}}A_0\eta - A_0\eta = 2A_0\eta - A_0\eta = A_0\eta = -A_0(2P_{\mathcal{S}} - 1)\eta.$$

Given a closed subspace $\mathcal{S} \subset \mathcal{H}_0$ such that $A_0(\mathcal{S}) \subset \mathcal{S}^\perp$ and $A_0(\mathcal{S}^\perp) \subset \mathcal{S}$, the orthogonal projection $E = P_{\mathcal{S}}$ satisfies that $EA_0E = (1 - E)A_0(1 - E) = 0$, and conversely. $\square$

Remark 2.3. Since A_0 is selfadjoint with trivial nullspace, the isometric part J_0, in the polar decomposition $A_0 = J_0|A_0^2| = |A_0|J_0$, is a symmetry. Note that a symmetry V anti-commutes with A_0, if and only if it anti-commutes with J_0. Indeed, V commutes with A_0^2 and with $|A_0|$, which has also trivial nullspace:

$$|A_0|J_0 V = -V|A_0|J_0 = -|A_0|VJ_0,$$

which implies that $J_0 V = -VJ_0$. The converse is trivial. Therefore, the following equivalent condition could be added in the previous theorem:

- Symmetries V in $\mathcal{H}_0$ such that $VJ_0 = -J_0 V$.

4

Remark 2.4. Note that $A \in \mathcal{D}$ is a selfadjoint contraction. Moreover, the existence of a symmetry intertwining A_0 with $-A_0$, means that the spectrum of the whole A is symmetric with respect to the origin, except for an eventual asymmetry at $\lambda = \pm 1$. For instance, if $0 < \lambda < 1$ is an eigenvalue of A, then $\pm\lambda \in \sigma(A_0)$, with the same multiplicity. The symmetry may break at $\lambda = 1$.

Remark 2.5. Consider now the following question: among the pairs $(P, Q) \in \mathcal{D}_A$, does the exist an optimal element which minimizes $P + Q$? We use the decomposition $\mathcal{H} = N(A) \oplus N(A^2 - 1) \oplus \mathcal{H}_0$ which reduces all pairs in $\mathcal{D}_A$. In the first subspace $N(A)$, all pairs are of the form (E, E), for E a projection onto a subspace of $N(A)$. Clearly, there is a minimal pair here, taking $E = 0$. On $N(A^2 - 1)$, there is one pair, and for this pair $P + Q$ equals the identity of $N(A^2 - 1)$. Let us prove that pairs (P_0, Q_0) in the generic part $\mathcal{H}_0$ are not comparable (unless they are equal). This implies that in the nontrivial case, where the generic part $\mathcal{H}_0 \neq 0$, there are no possible minimizers for $P + Q$. For $(P_0, Q_0) \in \mathcal{D}_{A_0}$, let V_0 be the corresponding Davis symmetry: $P_{V_0} = P_0$, $Q_{V_0} = Q_0$. Then

$$P_0 + Q_0 = 1 + V_0(1 - A_0^2)^{1/2}.$$

Thus, comparison of these sums is equivalent to comparison of the operators $V_0(1 - A_0^2)^{1/2} = (1 - A_0^2)^{1/2}V_0$. If V_1 is the symmetry corresponding to another pair (P_1, Q_1), then, since

$$\langle V_0(1 - A_0^2)^{1/2}\xi, \xi\rangle = \langle V_0(1 - A_0^2)^{1/4}\xi, (1 - A_0^2)^{1/4}\xi\rangle,$$

it follows that $P_1 + Q_1 \leq P_0 + Q_0$ if and only if

$$\langle V_0(1 - A_0^2)^{1/4}\xi, (1 - A_0^2)^{1/4}\xi\rangle \leq \langle V_1(1 - A_0^2)^{1/4}\xi, (1 - A_0^2)^{1/4}\xi\rangle.$$

Moreover, since $1 - A_0^2$ has trivial nullspace, $(1 - A_0^2)^{1/4}$ has dense range. Thus, the inequality above is equivalent to $V_1 \leq V_0$. This inequality is equivalent, in turn, to the inclusion $\mathcal{S}_1^+ \subset \mathcal{S}_0^+$, where $\mathcal{S}_i^+ = \{\xi \in \mathcal{H}_0 : V_i\xi = \xi\}$. Therefore our assumption $P_1 + Q_1 < P_0 + Q_0$ implies the existence of a nontrivial vector $\xi_0 \in \mathcal{S}_0^+$ such that $\xi_0 \perp \mathcal{S}_1^+$. This leads us to a contradiction. In fact, note that $A_0 V_0 = -V_0 A_0$ implies that $A_0(\mathcal{S}_i^+) \subset (\mathcal{S}_i^+)^\perp$ and $A_0((\mathcal{S}_i^+)^\perp) \subset \mathcal{S}_i^+$. Then

$$A_0\xi_0 \in (\mathcal{S}_0^+)^\perp \quad \text{and} \quad A_0\xi_0 \in A_0((\mathcal{S}_1^+)^\perp) \in \mathcal{S}_0^+,$$

i.e., $A_0\xi_0 = 0$, a contradiction, since A_0 has trivial nullspace.

2.1 Halmos decomposition

Given two projections P, Q, Halmos proved in [21] that there exists an isometric isomorphism between the generic part $\mathcal{H}_0$ (of P and Q) and a product space $\mathcal{L} \times \mathcal{L}$ which carries P_0 and Q_0 to the operator matrices

$$\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} C^2 & CS \\ CS & S^2 \end{pmatrix},$$

respectively. Here $C = \cos(\Gamma)$ and $S = \sin(\Gamma)$, where $0 \leq \Gamma \leq \pi/2$ is a positive operator in $\mathcal{L}$ with trivial nullspace. In particular, $CS = SC$ and S has trivial nullspace. Note that also C has trivial nullspace (i.e., $\pi/2$ is not an eigenvalue of X). Indeed, if $C\xi = 0$, then

$$\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}\begin{pmatrix} \xi \\ 0 \end{pmatrix} = \begin{pmatrix} \xi \\ 0 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} C^2 & CS \\ CS & S^2 \end{pmatrix}\begin{pmatrix} \xi \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix},$$

i.e., $\begin{pmatrix} \xi \\ 0 \end{pmatrix}$ lies in $R(P_0) \cap N(Q_0)$, which is trivial.

2.2 Friedrichs' angle

Given two closed subspaces $\mathcal{M}, \mathcal{N} \subset \mathcal{H}$, the Friedrichs angle [20] between $\mathcal{M}$ and $\mathcal{N}$ is the angle whose cosine is

$$c(\mathcal{M}, \mathcal{N}) := \sup\{|\langle \mu, \nu \rangle| : \mu \in \mathcal{M} \ominus \mathcal{N}, \nu \in \mathcal{N} \ominus \mathcal{M}, \|\mu\| = \|\nu\| = 1\}.$$

It holds that $c(\mathcal{M}, \mathcal{N}) = \|P_{\mathcal{M}} P_{\mathcal{N}} - P_{\mathcal{M} \cap \mathcal{N}})\|$ (see [14]).

We prove next that if $(P, Q) \in \mathcal{D}_A$, then $c(R(P), R(Q))$ does not depend on the pair, i.e., it is an invariant of A.

Proposition 2.6. *Friedrichs' angle* $c(R(P), R(Q)) = c(N(P), N(Q))$ *is constant for all* (P, Q) *in* $\mathcal{D}_A$.

Proof. Pick $(P', Q'), (P, Q) \in \mathcal{D}_A$. Let us reduce $P'Q' - P_{R(P') \cap R(Q')}$ and $PQ - P_{R(P) \cap R(Q)}$ in the three space decomposition (1). Note that $R(P) \cap R(Q)$ and $R(P') \cap R(Q')$ are non trivial only in $N(A)$. In $N(A)$ and $N(A^2 - 1) = R(P) \cap N(Q) \oplus N(P) \cap R(Q)$, $PQ - P_{R(P) \cap R(Q)}$ is trivial, and similarly for (P', Q'). In the generic part $\mathcal{H}_0$, by Theorem 3.1, there exists a unitary operator U such that $UP_0 U^* = P_0'$ and $UQ_0 U^* = Q_0'$. Then $U(R(P_0)) = R(P_0')$ and $U(R(Q_0)) = R(Q_0')$, so that

$$U(R(P_0) \cap R(Q_0)) = R(P_0') \cap R(Q_0') \ , \quad \text{i.e., } UP_{R(P_0) \cap R(Q_0)} U^* = P_{R(P_0') \cap R(Q_0')}.$$

Then, in the three space decomposition (1)

$$P'Q' - P_{R(P') \cap R(Q')} = 0 \oplus 0 \oplus U \left(P_0 Q_0 - P_{R(P_0) \cap R(Q_0)} \right) U^*,$$

and, thus,

$$\|P'Q' - P_{R(P') \cap R(Q')}\| = \|P_0 Q_0 - P_{R(P_0) \cap R(Q_0)}\| = \|PQ - P_{R(P) \cap R(Q)}\|,$$

i.e., $c(R(P'), R(Q')) = c(R(P), R(Q))$. $\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad\qquad \square$

Remark 2.7. Note that $c(R(P), R(Q)) = \|P_0 Q_0\|$. Using Halmos representation

$$\|P_0 Q_0\| = \|P_0 Q_0 P_0\|^{1/2} = \left\| \begin{pmatrix} C^2 & 0 \\ 0 & 0 \end{pmatrix} \right\|^{1/2} = \|C\|.$$

Then, the angle equals $\cos^{-1}(\|\cos(\Gamma)\|)$. If Γ is non invertible, $0 \in \sigma(\Gamma)$ and therefore $1 \in \sigma(C)$, and thus the angle is $\pi/2$. If Γ is invertible (which is equivalent to $A^2 - 1$ being of closed range), then $\|\cos(\Gamma)\| = \cos(\|\Gamma^{-1}\|^{-1})$, and the angle is $\|\Gamma^{-1}\|^{-1}$, or, equivalently, the lowest value in the spectrum of Γ. In any case, this quantity is an invariant of A. We shall see below (Remark 3.7), that $\|P_0 + Q_0\|$ is also an invariant of A_0.

3   A unitary action on $\mathcal{D}_A$

Let $\mathcal{A} := \{A\}' = \{T \in \mathcal{B}(\mathcal{H}) : TA = AT\}$. Since A is selfadjoint, $\mathcal{A}$ is a von Neumann subalgebra of $\mathcal{B}(\mathcal{H})$. Let $\mathcal{U}_{\mathcal{A}}$ be the unitary group of $\mathcal{A}$. Observe that $\mathcal{U}_{\mathcal{A}}$ is connected. The group $\mathcal{U}_{\mathcal{A}}$ acts on $\mathcal{D}_A$:

$$U \cdot (P, Q) = (UPU^*, UQU^*) \ , \ U \in \mathcal{U}_{\mathcal{A}} \ , \ (P, Q) \in \mathcal{D}_A,$$

because $UPU^* - UQU^* = UAU^* = A$.

The algebra $\mathcal{A}$ and the action of $\mathcal{U}_\mathcal{A}$ can be factored using the three space decomposition (1) $\mathcal{H} = N(A) \oplus N(A^2 - 1) \oplus \mathcal{H}_0$. The algebra $\mathcal{A}$ splits as

$$\mathcal{A} = \mathcal{B}(N(A)) \oplus \mathcal{A}_1 \oplus \mathcal{A}_0.$$

Let us describe the summands. The first summand is clearly $\mathcal{B}(N(A))$: any operator acting on $N(A)$ (and trivial in $N(A)^\perp$) commutes with A ($= 0$ in $N(A)$). A pair $(P, Q) \in \mathcal{D}_A$ reduces to (P', P') in $N(A)$, for some projection $P' \in \mathcal{B}(N(A))$. The action of the unitary group of $N(A)$ on these pairs is essentially the action of the unitary group of a space on the projections of the space. The orbits are parametrized by the dimensions of the range and the nullspace.

The second summand consists of the algebra of operators which commute with $A|_{N(A^2-1)}$. In the decomposition $N(A^2 - 1) = N(A - 1) \oplus N(A + 1)$, $A|_{N(A^2-1)}$ is the matrix

$$A|_{N(A^2-1)} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

and the operators in $N(A^2 - 1)$ which commute with A are of the form

$$T = \begin{pmatrix} T_1 & 0 \\ 0 & T_2 \end{pmatrix}.$$

The unitary operators of this form leave $A|_{N(A^2-1)}$ fixed (a fact consistent with the observation that all pairs $(P, Q) \in \mathcal{D}_A$ reduce to a unique element in $N(A^2 - 1)$).

The third summand is $\mathcal{A}_0 := \{A|_{\mathcal{H}_0} : A \in \mathcal{A}\}$. Therefore, it is natural to focus on the action of $\mathcal{U}_{\mathcal{A}_0}$, the unitary group of the part $\mathcal{A}_0$.

Theorem 3.1. *The action of $\mathcal{U}_{\mathcal{A}_0}$ on $\mathcal{D}_{\mathcal{A}_0}$ is transitive.*

Proof. Let $(P_0, Q_0), (P_0', Q_0') \in \mathcal{D}_{\mathcal{A}_0}$. Denote by V and V' the symmetries (which anti-commute with A_0) which correspond to these pairs. Consider the decomposition

$$\mathcal{H}_0 = N(P_0 + Q_0' - 1) \oplus N(P_0 + Q_0' - 1)^\perp.$$

Note that $N(P_0 + Q_0' - 1)$ reduces simultaneously both pairs $(P_0, Q_0), (P_0', Q_0')$. First note that $N(P_0 + Q_0' - 1) = N(P_0' + Q_0 - 1)$, because $P_0 - Q_0 = P_0' - Q_0'$. Also note that

$$N(P_0 + Q_0' - 1) = N(P_0 - (1 - Q_0')) = R(P_0) \cap R(1 - Q_0') \oplus N(P_0) \cap N(1 - Q_0')$$

$$= R(P_0) \cap N(Q_0') \oplus N(P_0) \cap R(Q_0'),$$

which reduces P_0 and Q_0', and similarly for P_0' and Q_0.

In the second subspace $N(P_0 + Q_0' - 1)^\perp$, the operator $P_0 + Q_0' - 1$ is selfadjoint and has trivial nullspace; therefore, in the polar decomposition

$$P_0 + Q_0' - 1 = \Sigma |P_0 + Q_0' - 1| = |P_0 + Q_0' - 1| \Sigma,$$

the operator Σ is a symmetry which, by the same argument as in the proof of Theorem 2.2, satifies

$$\Sigma P_0|_{N(P_0+Q_0'-1)^\perp} \Sigma = Q_0'|_{N(P_0+Q_0'-1)^\perp} \quad \text{and} \quad \Sigma Q_0'|_{N(P_0+Q_0'-1)^\perp} \Sigma = P_0|_{N(P_0+Q_0'-1)^\perp}.$$

Then,
$$\Sigma A_0|_{N(P_0+Q_0'-1)^\perp}\Sigma = -A_0|_{N(P_0+Q_0'-1)^\perp}.$$

The fact that $P_0 + Q_0' - 1 = P_0' + Q_0 - 1$, implies that this operator Σ intertwines also the reductions of P_0' and Q_0 to $N(P_0 + Q_0' - 1)^\perp$

The symmetry V, which is obtained (by means of the Borel functional calculus) as the sign function of $P_0 + Q_0 - 1$, also is reduced by $N(P_0 + Q_0' - 1)$. Clearly, $V|_{N(P_0+Q_0'-1)^\perp}$ also anticommutes with $A_0|_{N(P_0+Q_0'-1)^\perp}$. Then, the unitary operator U_1 in $N(P_0 + Q_0' - 1)^\perp$ defined as
$$U_1 = \Sigma V|_{N(P_0+Q_0'-1)^\perp},$$

commutes with $A_0|_{N(P_0+Q_0'-1)^\perp}$. Moreover, it satisfies

$$U_1 P_0|_{N(P_0+Q_0'-1)^\perp} U_1^* = \Sigma(V P_0 V)|_{N(P_0+Q_0'-1)^\perp}\Sigma = \Sigma Q_0|_{N(P_0+Q_0'-1)^\perp}\Sigma = P_0'|_{N(P_0+Q_0'-1)^\perp}$$

and, similarly,
$$U_1 Q_0|_{N(P_0+Q_0'-1)^\perp} U_1^* = Q_0'|_{N(P_0+Q_0'-1)^\perp}.$$

Let us find a unitary in the other subspace, $N(P_0 + Q_0' - 1)$. Trivially, $P_0|_{N(P_0+Q_0'-1)} = 1 - Q_0'|_{N(P_0+Q_0'-1)}$, and also $P_0'|_{N(P_0+Q_0'-1)} = 1 - Q_0'|_{N(P_0+Q_0'-1)}$. Note that

$$A_0 P_0|_{N(P_0+Q_0'-1)} = (1-Q_0)P_0|_{N(P_0+Q_0'-1)} = P_0'(1-Q_0')|_{N(P_0+Q_0'-1)} = P_0'A_0|_{N(P_0+Q_0'-1)},$$

and, similarly,
$$A_0 Q_0|_{N(P_0+Q_0'-1)} = Q_0'A_0|_{N(P_0+Q_0'-1)}.$$

Then, again by the same argument as above (and as in the proof of Theorem 2.2), the isometric part U_2 in the polar decomposition of $A_0|_{N(P_0+Q_0'-1)}$ (which has trivial nullspace in the whole $\mathcal{H}_0$ and, thus also in $N(P_0 + Q_0' - 1)$), is a symmetry ($U_2 = U_2^*$) which satisfies

$$U_2 P_0|_{N(P_0+Q_0'-1)} U_2 = P_0'|_{N(P_0+Q_0'-1)} \quad \text{and} \quad U_2 Q_0|_{N(P_0+Q_0'-1)} U_2 = Q_0'|_{N(P_0+Q_0'-1)}$$

In particular, this implies that U_2 commutes with $A_0|_{N(P_0+Q_0'-1)}$. Consider, then,

$$U = U_2 \oplus U_1 \quad \text{acting in} \quad N(P_0 + Q_0' - 1) \oplus N(P_0 + Q_0' - 1)^\perp = \mathcal{H}_0.$$

Clearly, U is a unitary operator which commutes with A_0, and satisfies $U \cdot (P_0, Q_0) = (P_0', Q_0')$. □

The following result appeared in [30]. It was proved there using a different technique: Shi, Ji and Du obtained a parametrization of $\mathcal{D}_{A_0}$, in terms of unitaries in a von Neumann algebra. The result is proved here as an easy consequence of the above theorem:

Corollary 3.2. $\mathcal{D}_{A_0}$ *is connected. The connected components of $\mathcal{D}_A$ are parametrized by the connected components of the space of projections $\mathcal{P}(N(A))$ of the space $N(A)$.*

Proof. $\mathcal{U}_{A_0}$ is the unitary group of a von Neumann algebra, therefore connected, and the map $\pi_{(P_0,Q_0)}$ is continuous. The assertion on the components of $\mathcal{D}_A$ follows form the description of $\mathcal{D}_A$ done in Section 1. □

Remark 3.3. The unitary U obtained above is, in fact, an explicit formula in terms of P_0, Q_0, P_0' and Q_0'. However, if one fixes for instance the pair (P_0, Q_0), U is not a continuous formula in terms of (P_0', Q_0') (a continuous formula would provide a continuous global cross-section for the action). Indeed, the formula of U depends on the decomposition $\mathcal{H}_0 = N(P_0 + Q_0' - 1) \oplus N(P_0 + Q_0' - 1)^\perp$. Or, equivalently, on the map

$$Q_0' \mapsto P_{N(P_0 + Q_0' - 1)} = P_{N(P_0 - (1 - Q_0'))}.$$

One can find trivial examples (in dimension 2, for instance) where this map is not continuous.

However, in some cases the action does have continuous *local* cross sections. Let us show one such case.

Given a fixed $(P_0, Q_0) \in \mathcal{D}_{A_0}$, consider the continuous (surjective) map

$$\pi_{(P_0, Q_0)} : \mathcal{U}_{A_0} \to \mathcal{D}_{A_0} , \quad \pi_{(P_0, Q_0)}(U) = U \cdot (P_0, Q_0) = (U P_0 U^*, U Q_0 U^*).$$

Lemma 3.4. $1 - A^2$ *has closed range if and only if for any* $(P_0, Q_0) \in \mathcal{D}_{A_0}$, $P_0 + Q_0 - 1$ *is invertible.*

Proof. Suppose that $1 - A^2$ has closed range. Note the formula (see [23] p. 33, or compute directly):

$$(P - Q)^2 + (P + Q - 1)^2 = 1,$$

or, equivalently, $1 - A^2 = (P + Q - 1)^2$. It follows that $(P + Q - 1)^2$ has closed range. In the generic part $\mathcal{H}_0$, $(P + Q - 1)^2|_{\mathcal{H}_0} = (P_0 + Q_0 - 1)^2$ has trivial nullspace. Thus, $(P_0 + Q_0 - 1)^2$ is invertible, and, thus, also $P_0 + Q_0 - 1$ is invertible.

Conversely, if $P_0 + Q_0 - 1$ is invertible, then $(P_0 + Q_0 - 1)^2$ is also invertible, and then its extension $(P + Q - 1)^2 = 1 - A^2$ (which is zero in $N(P + Q - 1)$) has closed range. $\square$

Remark 3.5. Using Halmos decomposition, a simple computation shows that $R(A^2 - 1)$ is closed, which means that $A_0^2 - 1$ is invertible if and only if S (or Γ) is invertible in $\mathcal{L}$.

For such A as above, the map $\pi_{(P_0, Q_0)}$ has continuous *local* cross-sections.

Proposition 3.6. *Let* $A \in \mathcal{D}$ *such that* $A^2 - 1$ *has closed range. Then the map* $\pi_{(P_0, Q_0)}$ *has continuous local cross-sections.*

Proof. Consider the set

$$\{(P', Q') \in \mathcal{D}_{A_0} : P_0 + Q' - 1 \text{ is invertible in } \mathcal{H}_0\}.$$

Since the set of invertible operators is open, this set is clearly an open subset of $\mathcal{D}_{A_0}$ (considered with the relative topology of $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$). It is a neighbourhood of (P_0, Q_0): if $(P', Q') = (P_0, Q_0)$, $P_0 + Q_0 - 1$ is invertible, by the above Lemma. Then the map

$$\mathbf{s} : \{(P', Q') \in \mathcal{D}_{A_0} : P_0 + Q' - 1 \text{ is invertible in } \mathcal{H}_0\} \to \mathcal{U}_{A_0}, \ \mathbf{s}(P', Q') = sgn(P_0 + Q' - 1)V$$

is continuous. Here $sgn(P_0 + Q' - 1)$ denotes the sign of the selfadjoint (invertible) operator $P_0 + Q' - 1$, $sgn(P_0 + Q' - 1) = (P_0 + Q' - 1)\left((P_0 + Q' - 1)^2\right)^{-1/2}$. Note that the function sgn is continuous on the set of invertible operators. As seen above, it is an element of $\mathcal{U}_{A_0}$ (called ΣV in the proof of Theorem 3.1). Also, it is clear that it is a cross section in a neighbouthood of (P_0, Q_0). One obtains cross sections around other points by translating this map using the transitive action. $\square$

Remark 3.7. In Remark 2.5, we observed that if $(P_0, Q_0), (P_0', Q_0')$ are the generic parts of two pairs $(P, Q), (P', Q') \in \mathcal{D}_A$, the operators $P_0 + Q_0$ and $P_0' + Q_0'$ are not comparable (an inequality implies equality). The transitivity of the action of $\mathcal{U}_{\mathcal{A}_0}$ on $\mathcal{D}_{A_0}$ implies that the norms of these operators coincide. Indeed, since there exists a unitary operator U in $\mathcal{A}_0$ such that $UP_0U^* = P_0'$ and $UQ_0U^* = Q_0'$, it follows that $U(P_0+Q_0)U^* = P_0'+Q_0'$, and therefore $\|P_0+Q_0\| = \|P_0'+Q_0'\|$.

4   A presentation of $\mathcal{A}_0$ in terms of Halmos decomposition

Proposition 4.1. *The algebra $\mathcal{A}_0$, represented in $\mathcal{B}(\mathcal{L} \times \mathcal{L})$, consists of matrices of the form*

$$\left\{ \begin{pmatrix} X & Y \\ Y & Z \end{pmatrix} : X, Y, Z \in \mathcal{B}(\mathcal{L}) \text{ commute with } \Gamma, \text{ and } C(X - Z) + 2SY = 0 \right\}.$$

Proof. In $\mathcal{L} \times \mathcal{L}$, A_0 is

$$\begin{pmatrix} S^2 & -CS \\ -CS & -S^2 \end{pmatrix}.$$

Let $\begin{pmatrix} X & Y_1 \\ Y_2 & Z \end{pmatrix}$ be an operator which commutes with A_0. Then, in particular, it commutes with A_0^2 which is given by

$$\begin{pmatrix} S^2 & 0 \\ 0 & S^2 \end{pmatrix}.$$

Then X, Y_1, Y_2, Z commute with S^2. Therefore, they commute also with its square root $|S| = S$, and with C. Thus, X, Y_1, Y_2, Z commute with $e^{i\Gamma} = C + iS$, and with its analytic logarithm $i\Gamma$ (since $\|\Gamma\| \le \pi/2 < \pi$). Straightforward computations show that an operator lies in the commutant of A_0 if and only if

- $CSY_1 = CSY_2$, which means that $Y_1 = Y_2$, because C, S have trivial nullspaces, and

- $S^2Y - CSZ = -CSX - S^2Y$, which, again using that S has trivial nullspace, means that $C(X - Z) + 2SY = 0$.

$\square$

Remark 4.2. Note that since $S = \sin(\Gamma)$ has trivial nullspace, then the (eventually unbounded, densely defined) operator $\tau = \tan(\Gamma)$ is defined, and the condition $C(X - Z) + 2SY = 0$ can be replaced by

$$Z = X + 2\tau Y.$$

In particular, it implies that $\tau Y = Y\tau$ is bounded.

Remark 4.3. It is also easy to characterize the unitaries in $\mathcal{A}_0$ which leave (P_0, Q_0) fixed. They are the (unitary) matrices which commute with P_0 and Q_0. The first relation implies that they must be diagonal matrices. Commutation with the second projection implies, after simple computations (using that C has trivial nullspace), that they are of the form

$$\{W \in \mathcal{U}_{\mathcal{A}_0} : WP_0W^* = P_0 \text{ and } WQ_0W^*\} = \left\{ \begin{pmatrix} W' & 0 \\ 0 & W' \end{pmatrix} : W' \in \mathcal{U}(\mathcal{L}), W'\Gamma = \Gamma W' \right\}.$$

Using this representation of the generic part, we can further analize the condition that $A^2 - 1$ has closed range. Recall from Lemma 3.4, that this is equivalent to the invertibility of $P_0 + Q_0 - 1$, for any pair $(P_0, Q_0) \in \mathcal{D}_{A_0}$.

Proposition 4.4. *The following are equivalent*

1. $A^2 - 1$ *has closed range.*

2. $PQP - P$ *has closed range.*

3. $P_0 Q_0 P_0 - P_0$ *is invertible in* $R(P_0)$

4. Γ *is invertible in* $\mathcal{L}$.

Proof. Clearly, it suffices to examine the reductions to the generic part $\mathcal{H}_0$. Using Halmos representation, one gets

$$A_0^2 - 1 = \begin{pmatrix} -S^2 & 0 \\ 0 & -S^2 \end{pmatrix},$$

and $P_0 Q_0 P_0 - P_0 = -S^2 = -\sin(\Gamma)$. The equivalence of these conditions becomes apparent. $\square$

5 Examples

We present examples of operators A, which will be the object of further study. The first one has continuous spectrum.

Example 5.1. Let $\mathcal{H} = L^2(-1, 1)$ and $A = M_t$ (multiplication by the variable): $Af(t) = tf(t)$. Note that A anti-commutes with the symmetry V, $Vf(t) = f(-t)$. Therefore $A = P_V - Q_V$, following the notation of Davis' characterization in Section 1, and both projections can be computed explicitly. Since A has no eigenvalues, it follows that $\mathcal{H}_0 = \mathcal{H}$ (i.e., P_V, Q_V or any pair of projections with difference A are in generic position). Also note that the algebra $\mathcal{A}$ is $L^\infty(-1, 1)$, represented as multiplication operators in $\mathcal{H}$. Therefore, if one chooses to parametrize elements in $\mathcal{D}_A$ by means of isometries, $\mathcal{D}_A$ consists of all symmetries V_φ of the form

$$V_\varphi = M_\varphi V M_{\bar{\varphi}},$$

for $\varphi \in L^\infty(-1, 1)$, with $|\varphi(t)| = 1$ *a.e.*, modulo the commutant of V, i.e., the unimodular functions of $L^\infty(-1, 1)$ which are *essentially even*. Explicitely,

$$V_\varphi f(t) = \varphi(t)\bar{\varphi}(-t)f(t),$$

modulo the functions φ such that $\varphi(t) = \varphi(-t)$ for almost every $t \in (-1, 1)$.

The second example has pure point spectrum.

Example 5.2. Let $I, J \subset \mathbb{R}^n$ be Lebesgue measurable sets with positive finite measure. Consider $\mathcal{H} = L^2(\mathbb{R}^n)$ and the projections

$$P_I = M_{\chi_I} \quad \text{and} \quad Q_J = \mathcal{F}^{-1} P_J \mathcal{F}$$

where χ_D is the characteristic function of $D \subset \mathbb{R}^n$ and $\mathcal{F} : \mathcal{H} \to \mathcal{H}$ is the Fourier-Plancherel transform. These pairs have been studied in connection with mathematical formulation of the

uncertainty principle (see [25], the survey [19] or the book [22]). Specifically, the products $P_I Q_J$, $P_I Q_J P_I$ are of interest. Among the basic facts concerning these operators, it is known that they are Hilbert-Schmidt operators, and that

$$R(P_I) \cap R(Q_J) = R(P_I) \cap N(Q_J) = N(P_I) \cap R(Q_J) = \{0\},$$

and $N(P_I) \cap N(Q_J)$ is infinite dimensional (see for instance [25]). In particular, $P_I Q_J P_I$ has a complete orthonormal system of eigenvectors (i.e., $P_I Q_J P_I$ is diagonalizable). In [5] it was proved that for a pair of projections P, Q, PQP is diagonalizable if and only if $P - Q$ is diagonalizable. In this case, there is an explicit relation between the eigenvectors and eigenvalues of PQP and $P - Q$. If s_n are the eigenvalues of PQP ($0 < s_n < 1$), then $\pm\lambda_n = \pm(1 - s_n)^{1/2}$ are the eigenvalues of $P - Q$. The eigenvalue $s = 1$ corresponds with $\lambda = 0$.

Thus, our second example $A = P_I - Q_J$ is diagonalizable. Moreover,

$$\mathcal{H}_0 = N(A)^\perp = (N(P_I) \cap N(Q_J))^\perp .$$

In the particular case $I = (0, 1)$ and $J = (-\Omega/2, \Omega/2)$ the eigenvectors are known (called prolate spheroidal functions [31], [22]), and the eigenvalues have simple multiplicity. Therefore, in this case $\mathcal{A}_0$ consists of all diagonal matrices in this orthonormal basis. In particular, $\mathcal{A}_0$ is commutative, as in the previous example.

Let us characterize in this example the symmetries which anti-commute with A_0. Note that this implies that V commutes with A_0^2. Therefore V has block diagonal form, with blocks of size 2×2, generated, for each fixed $n \geq 1$, by the eigenvectors e_n, f_n of λ_n and $-\lambda_n$, respectively. Note that $V e_n$ is an eigenvector for $-\lambda_n$:

$$A_0 V e_n = -V A_0 e_n = -\lambda_n V e_n.$$

Thus, since in this case all eigenvalues have multiplicity one, $V e_n = \omega_n f_n$, for some $\omega_n \in \mathbb{C}$ with $|\omega_n| = 1$. Similarly, $V f_n$ is an unimodular multiple of e_n. The fact that $V^2 = 1$ implies that $V f_n = \bar{\omega}_n e_n$. Therefore, any symmetry V anti-commuting with A_0 (in $\mathcal{H}_0$) is of the form

$$V = V_\omega = \oplus_{n=1}^\infty V_{\omega_n} \in \oplus_{n=1}^\infty \mathcal{H}_n , \quad \text{where} \quad V_{\omega_n} = \begin{pmatrix} 0 & \omega_n \\ \bar{\omega}_n & 0 \end{pmatrix}$$

and $\mathcal{H}_n$ is the subspace spanned by e_n and f_n. That is, the elements of $\mathcal{D}_{A_0}$ can be parametrized by sequences $\omega = \{\omega_n\}$ of complex numbers of modulus 1.

Remark 5.3. If we consider the 2×2 matrix representation of the previous section, elementary computations show that if $\mathcal{A}_0$ is commutative, then all entries in the matrices commute.

Example 5.4. Consider $\mathcal{H} = L^2(\mathbb{T})$, put $\mathcal{H}^+ = H^2(\mathbb{T})$, $P_+ = P_{\mathcal{H}^+}$. Let $\mathbf{a} = \{a_1, \dots, a_N\}$, $\mathbf{b} = \{b_1, \dots, b_N\}$ two (finite) sequences of points in the open disk $\mathbb{D}$. We suppose that $a_i \neq a_j$ if $i \neq j$ and $a_i \neq b_j$ for all i, j. Let $B_\mathbf{a}, B_\mathbf{b}$ be the corresponding Blaschke products. Put $P_\mathbf{a} = P_{B_\mathbf{a} \mathcal{H}^+}$, and similarly $P_\mathbf{b} = P_{B_\mathbf{b} \mathcal{H}^+}$. Note that since the multiplication operator $M_{B_\mathbf{a}}$ is a unitary operator in $\mathcal{H}$, $P_\mathbf{a} = M_{B_\mathbf{a}} P_+ M_{\bar{B}_\mathbf{a}}$ and similarly $P_\mathbf{b} = M_{B_\mathbf{b}} P_+ M_{\bar{B}_\mathbf{b}}$. Also note that, if $\mathbf{a}\#\mathbf{b} = \{a_1, \dots, a_N, b_1, \dots, b_N\}$, then

$$\mathcal{H}_\mathbf{a} \cap \mathcal{H}_\mathbf{b} = \mathcal{H}_{\mathbf{a}\#\mathbf{b}} ;$$

$$\mathcal{H}_\mathbf{a} \cap \mathcal{H}_\mathbf{b}^\perp = \mathcal{H}_\mathbf{a}^\perp \cap \mathcal{H}_\mathbf{b} = \{0\}.$$

This is an exercise which follows from the fact that $\mathbf{a}$ and $\mathbf{b}$ have the same cardinality (see also [3] for a proof). Moreover,

$$\mathcal{H}_{\mathbf{a}}^{\perp} \cap \mathcal{H}_{\mathbf{b}}^{\perp} = (\mathcal{H}_{\mathbf{a}} \cup \mathcal{H}_{\mathbf{b}})^{\perp} = (\mathcal{H}^+)^{\perp} = \mathcal{H}^-$$

because $B_{\mathbf{a}}$ and $B_{\mathbf{b}}$ are coprime inner functions ($a_i \neq b_j$). Therefore, the generic part $\mathcal{H}_0$ equals the model space

$$\mathcal{H}_0 = \mathcal{H}^+ \ominus \mathcal{H}_{\mathbf{a}\#\mathbf{b}},$$

spanned by $\{k_{a_1}, \ldots, k_{a_N}, k_{b_1}, \ldots, k_{b_N}\}$, where k_c denotes the reproducing kernel of $\mathcal{H}^+$ at $c \in \mathbb{D}$. Since we have chosen different points a_i, b_j in the disk, these functions are linearly independent, and $\mathcal{H}_0$ has dimension $2N$. Thus $\mathcal{A}_0$ is a finite dimensional algebra. The operator $A_0 = P_{\mathbf{a}} - P_{\mathbf{b}}|_{\mathcal{H}_0}$ is a $2N \times 2N$ matrix. Let us compute A_0 in the elements of the basis k_{a_i}, k_{b_j}. First, note that $k_{a_i} \in \mathcal{H}_a^{\perp}$: if $h \in \mathcal{H}^+$, $\langle B_{\mathbf{a}}h, k_{a_i} \rangle = B_{\mathbf{a}}(a_i)h(a_i) = 0$. Then

$$P_{\mathbf{a}}k_{a_i} = 0 \ , \quad P_{\mathbf{b}}k_{b_j} = 0.$$

Note that

$$P_+ \bar{B}_{\mathbf{a}} k_{b_j} = \sum_{l \geq 0} \langle \bar{B}_{\mathbf{a}} k_{b_j}, z^l \rangle z^l = \sum_{l \geq 0} \langle k_{b_j}, B_{\mathbf{a}} z^l \rangle z^l = \sum_{l \geq 0} \overline{B_{\mathbf{a}}(b_j)} (\bar{b}_j z)^l = \overline{B_{\mathbf{a}}(b_j)} \frac{1}{1 - \bar{b}_j z} = \overline{B_{\mathbf{a}}(b_j)} k_{b_j}.$$

Then, $P_{\mathbf{a}} k_{b_j} = M_{B_{\mathbf{a}}} P_+ M_{\bar{B}_{\mathbf{a}}} k_{b_j} = \overline{B_{\mathbf{a}}(b_j)} B_{\mathbf{a}} k_{b_j}$. Similarly for $P_{\mathbf{b}} k_{a_i}$. Then,

$$A_0 k_{a_i} = -\overline{B_{\mathbf{b}}(a_i)} B_{\mathbf{b}} k_{a_i}$$

and

$$A_0 k_{b_j} = \overline{B_{\mathbf{a}}(b_j)} B_{\mathbf{a}} k_{b_j}.$$

Thus, in principle, it is possible to compute the $2N \times 2N$ matrix of A_0 in the (non-orthogonal) basis of the reproducing kernels k_{a_i}, k_{b_j}. It would be interesting to know if under the present assumptions, A_0 has eigenvalues of simple multiplicity.

Example 5.5. Let $E \in \mathcal{B}(\mathcal{H})$ be a non-selfadjoint idempotent operator ($E^2 = E$). Consider the orthogonal projections $P_{R(E)}$ and $P_{R(E^*)} = P_{N(E)^{\perp}}$. In matrix form, in terms of the decomposition $\mathcal{H} = R(E) \oplus R(E)^{\perp}$, E is written

$$E = \begin{pmatrix} 1 & B \\ 0 & 0 \end{pmatrix},$$

where $B : R(E)^{\perp} \to R(E)$. Consider the selfadjoint operator $S = E + E^* - 1$. S is selfadjoint with trivial nullspace, which satisfies $SE = E^*S$ and $SE^* = ES$. Then, similarly as before, the unitary part of S in the polar decomposition intertwines E and E^*. Then, it also intertwines the range projections. Moreover, by straightforward matrix computations (which were done explicitly in [6]), this unitary part coincides with Davis' symmetry V for the pair of projections $P_{R(E)}, P_{R(E^*)}$. Also note the well known formulas

$$P_{R(E)} = ES^{-1} \quad \text{and} \quad P_{R(E^*)} = E^*S^{-1},$$

so that

$$A = (E - E^*)S^{-1}. \tag{2}$$

Note that $R(E) \cap N(E) = \{0\}$ and $R(E^*) \cap N(E^*) = \{0\}$. Straightforward computation show that

$$R(E) \cap R(E^*) = N(B^*) \quad \text{and} \quad N(E^*) \cap N(E) = N(B).$$

Thus, in order that $P_{R(E)}$ and $P_{R(E^*)}$ be in generic position, B should have dense range and trivial nullspace. Let us assume this. In particular it implies that $\dim R(E) = \dim R(E)^\perp$. Clearly, if we want to study the structure of $\mathcal{D}_A$, we can replace E with UEU^*, where $U : \mathcal{H} \to \mathcal{J}$ is a unitary transformation. Thus, $A = P_{R(E)} - P_{R(E^*)}$ is replaced by UAU^*. Pairs in $\mathcal{D}_A$ are mapped onto pairs in $\mathcal{D}_{UAU^*}$ by means of $(P, Q) \mapsto (UPU^*, UQU^*)$. Therefore (by the equality of dimensions between $R(E)$ and $R(E)^\perp$) we can choose a model $\mathcal{J} = \mathcal{L} \times \mathcal{L}$. Next consider the polar decomposition of B, $B = W_0|B|$. Clearly $W_0 : \mathcal{L} \to \mathcal{L}$ is a unitary operator. Consider the unitary operator W in $\mathcal{L} \times \mathcal{L}$ given by $W = \begin{pmatrix} W_0^* & 0 \\ 0 & 1 \end{pmatrix}$. Then,

$$WEW^* = \begin{pmatrix} 1 & W_0^* B \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 1 & |B| \\ 0 & 0 \end{pmatrix}.$$

Summarizing, we can suppose that $\mathcal{H} = \mathcal{L} \times \mathcal{L}$ and B is positive with trivial nullspace. Therefore, with the current assumptions, using (2), one has

$$A_0 = A = \begin{pmatrix} B^2(1 + B^2)^{-1} & -B(1 + B^2)^{-1} \\ -B(1 + B^2)^{-1} & -B^2(1 + B^2)^{-1} \end{pmatrix}.$$

In order to describe $\mathcal{A}_0 = \{A_0\}'$, note that $A_0^2 = \begin{pmatrix} B^2(1 + B^2)^{-1} & 0 \\ 0 & B^2(1 + B^2)^{-1} \end{pmatrix}$. Therefore, if $X = \begin{pmatrix} X_{11} & X_{12} \\ X_{21} & X_{22} \end{pmatrix}$ belongs to $\mathcal{A}_0$, in particular it commutes with A_0^2. This clearly implies that the entries X_{ij} commute with B (recall that $B \geq 0$). Next, note that the condition that X commutes with A_0 means that

$$BX_{12} = X_{21}B = BX_{21},$$

which implies $X_{12} = X_{21}$, because $N(B) = 0$, and that

$$BX_{22} - X_{11}B = B(X_{22} - X_{11}) = 2B^2 X_{12},$$

which implies that $X_{22} = 2BX_{12} + X_{11}$. Therefore

$$\mathcal{A}_0 = \{ \begin{pmatrix} Y & Z \\ Z & Y + 2BZ \end{pmatrix} : Y, Z \text{ commute with } B \}.$$

Let us describe the isotropy subalgebra, i.e., the operators which commute with $P_{R(E)}$ and $P_{R(E^*)}$. Operators which commute with $P_{R(E)} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$, are diagonal matrices. If they belong additionally to $\mathcal{A}_0$, they are of the form

$$\{ \begin{pmatrix} Y & 0 \\ 0 & Y \end{pmatrix} : Y \text{ commutes with } B \}.$$

Easy examples (of positive operators B), show that the isotropy subalgebra, and therefore $\mathcal{A}_0$, may not be commutative.

6 A regular structure for D_A

We shall prove that $\mathcal{D}_A$ is an homogeneous C^∞ space of the unitary group $\mathcal{U}_A$. If $A^2 - 1$ has closed range, then $\mathcal{D}_A$ is additionally a complemented submanifold of $\mathcal{P}(\mathcal{H}) \times \mathcal{P}(\mathcal{H})$. As seen above, $\mathcal{D}_A$ decomposes as three spaces in the decomposition (1).

- In $N(A)$, the group acting is the whole unitary group of $N(A)$, and the space $\mathcal{D}_A$ reduces to pairs of the form (E, E), where $E \in \mathcal{P}(\mathcal{H})$ and $R(E) \subset N(A)$, i.e., $\mathcal{D}_A|_{N(A)}$ identifies with the space of projections in the Hilbert space $N(A)$, under the action of the unitary group of $N(A)$. This space is well studied: it is a C^∞ complemented submanifold of $\mathcal{B}(N(A))$ (see [12]).

- In $N(A^2 - 1)$, $\mathcal{D}_A|_{N(A^2-1)}$ is a single point , namely, $(P_{N(A-1)}, P_{N(A+1)})$.

- Therefore, the task is reduced to show that $\mathcal{D}_{A_0}$ has local regular structure.

In order to prove that $\mathcal{D}_{A_0}$ has differentiable structure, and also in order to define later a linear connection in this manifold, the following map will be useful:

Definition 6.1. *Fix* $(P_0, Q_0) \in \mathcal{D}_{A_0}$, *and fix also a Halmos decomposition for this pair. Consider the map* $E_{(P_0,Q_0)} : \mathcal{A}_0 \to \mathcal{A}_0$,

$$E_{(P_0,Q_0)}\left(\begin{pmatrix} X & Y \\ Y & Z \end{pmatrix} \right) = \begin{pmatrix} \frac{1}{2}(X + Z) & 0 \\ 0 & \frac{1}{2}(X + Z) \end{pmatrix}. \tag{3}$$

It is easy to see that this map is a conditional expectation, with range equal to the subalgebra of elements in $\mathcal{A}_0$ which commute with P_0 and Q_0 (see Remark 4.3).

Remark 6.2. Let us prove that the conditional expectation $E_{(P_0,Q_0)}$ depends only on the pair (P_0, Q_0) (and not on the Halmos decomposition). Indeed, first note that $Q_0 - P_0 Q_0 P_0 - P_0^\perp Q_0 P_0^\perp = \begin{pmatrix} 0 & CS \\ CS & 0 \end{pmatrix}$. Let us denote this operator by K. Since C and S have trivial nullspace, then $N(K) = 0$. Also, it is clear that $K^* = K$. Then, in the polar decomposition of K,

$$K = W|K| = W \begin{pmatrix} CS & 0 \\ 0 & CS \end{pmatrix};$$

using again that C and S have trivial nullspaces, it follows that it must be $W = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. The operator W is obtained by means of the functional calculus of K: $W = sgn(K)$, where sgn denotes the (Borel, eventually non continuous) sign function ($sgn(t) = 1$ if $t \geq 0$, -1 if $t < 0$). Note that $P_0 W = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$ and $WP_0 = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$. Then, if $M = \begin{pmatrix} X & Y \\ Y & Z \end{pmatrix}$, we get

$$E_{(P_0,Q_0)}(M) = \frac{1}{2} P_0 (M + WMW) P_0 + \frac{1}{2} P_0^\perp (M + WMW) P_0^\perp.$$

Definition 6.3. *As above, fix* $(P_0, Q_0) \in \mathcal{D}_{A_0}$ *and a Halmos decomposition for this pair. Denote*

$$\mathbb{H}_{(P_0,Q_0)} := N(E_{(P_0,Q_0)} \cap (\mathcal{A}_0)_{ah} = \{ \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix} : Y^* = -Y, Y\Gamma = \Gamma Y, Y\tau \text{ is bounded in } \mathcal{L} \}. \tag{4}$$

Note that if $Y\tau$ *is bounded, then* $Z - X = 2Y\tau$.

Remark 6.4. Let us point out the fact, which will be relevant later, that for any (matrix) element in $\mathbb{H}_{(P_0,Q_0)}$, all entries of the matrix commute. Indeed, Y commutes with $\tau = \tan(\Gamma)$.

In order to study the local structure of $\mathcal{D}_{A_0}$ we first suppose that $R(A^2-1)$ is closed. In this case we shall prove that $\mathcal{D}_{A_0}$ is a submanifold of $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$ (as well a a homogeneous space of $\mathcal{U}_{A_0}$). To prove this fact we shall need the following lemma, which is an application of the inverse function theorem in Banach spaces. One can find a detailed and elementary proof of this fact in [29].

Lemma 6.5. *Let G be a Banach-Lie group acting smoothly on a Banach space X. For a fixed $x_0 \in X$, denote by $\pi_{x_0} : G \to X$ the smooth map $\pi_{x_0}(g) = g \cdot x_0$. Suppose that:*

1. *π_{x_0} is an open mapping, regarded as a map from G onto the orbit $\{g \cdot x_0 : g \in G\}$ of x_0 (with the relative topology of X).*

2. *The differential $d(\pi_{x_0})_1 : (TG)_1 \to X$ splits: its nullspace and range are closed complemented subspaces in the Banach-Lie algebra $\mathcal{G}$ of G and X, respectively.*

Then, the orbit $\{g \cdot x_0 : g \in G\}$ is a smooth submanifold of X, and the map

$$\pi_{x_0} : G \to \{g \cdot x_0 : g \in G\}$$

is a smooth submersion.

Proposition 6.6. *Suppose that $A^2 - 1$ has closed range. Then, $\mathcal{D}_{A_0}$ is a complemented C^∞ submanifold of $\mathcal{B}(\mathcal{H}) \times \mathcal{B}(\mathcal{H})$, and for any fixed $(P_0, Q_0) \in \mathcal{D}_{A_0}$, the map*

$$\pi_{(P_0,Q_0)} : \mathcal{U}_{A_0} \to \mathcal{D}_{A_0} \ , \ \pi_{(P_0,Q_0)}(U) = (UP_0U^*, UQ_0U^*)$$

is a C^∞ submersion.

Proof. We shall apply Lemma 6.5 above. Note that the condition that $\pi_{(P_0,Q_0)}$ is open is fullfilled: if $A^2 - 1$ has closed range, then $\pi_{(P_0,Q_0)}$ has continuous local cross-sections. A cross section on a neighbourhood of (P_0, Q_0) was defined in Section 2 by

$\mathbf{s} : \{(P', Q') \in \mathcal{D}_{A_0} : P_0 + Q' - 1 \in Gl(\mathcal{H}_0)\} \to \mathcal{U}_{A_0} \ , \ \mathbf{s}(P', Q') = (P_0 + Q' - 1)|P_0 + Q' - 1|^{-1}V.$

This map can be extended to a map $\check{\mathbf{s}}$ defined on an open subset in $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$, with values in $Gl(\mathcal{H}_0)$. Namely

$\check{\mathbf{s}} : \{(T, S) \in \mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0) : P_0 + S - 1 \in Gl(\mathcal{H}_0)\} \ , \ \check{\mathbf{s}}(T, S) = (P_0 + S - 1)|P_0 + S - 1|^{-1}V.$

Clearly, $\{(T, S) \in \mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0) : P_0 + S - 1 \in Gl(\mathcal{H}_0)\}$ is an open subset of $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$ containing (P_0, Q_0), and $\check{\mathbf{s}}$ is C^∞.

The differential $d(\pi_{(P_0,Q_0)})_1 : (\mathcal{A}_0)_{ah} \to \mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$ is given by

$$d(\pi_{(P_0,Q_0)})_1(Z) = (ZP_0 - P_0Z, ZQ_0 - Q_0Z).$$

Here, $(\mathcal{A}_0)_{ah}$ denotes the set of anti-Hermitian elements of $\mathcal{A}_0$ (which is the Banach-Lie algebra of $\mathcal{U}_{A_0}$). Clearly, this map has a natural extension

$$\Pi : \mathcal{B}(\mathcal{H}_0) \to \mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0) \ , \ \Pi(X) = (XP_0 - P_0X, XQ_0 - Q_0X).$$

16

Denote $\check{\mathbf{S}} = d(\check{\mathbf{s}})_{(P_0,Q_0)}$. The fact that $\mathbf{s}$ is a cross section for $\pi_{(P_0,Q_0)}$ implies that $\pi_{(P_0,Q_0)} \circ \mathbf{s} \circ \pi_{(P_0,Q_0)} = \pi_{(P_0,Q_0)}$. Equivalently,

$$\pi_{(P_0,Q_0)} \circ \check{\mathbf{s}} \circ \pi_{(P_0,Q_0)} = \pi_{(P_0,Q_0)}.$$

This is a composition of C^∞ maps defined on open subsets of Banach spaces. If we differentiate this identity at 1, we get

$$\Pi \circ \check{\mathbf{S}} \circ \Pi = \Pi. \tag{5}$$

If we restrict this identity to $(\mathcal{A}_0)_{ah}$, the image $\Pi((\mathcal{A}_0)_{ah})$ equals the image of $d(\pi_{(P_0,Q_0)})_1$. Then, identity (5) above implies that $\Pi \circ \check{\mathbf{S}}$ is an idempotent whose range equals the range of Π. It follows that the range of $d(\pi_{(P_0,Q_0)})_1$ is complemented in $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$.

The nullspace of $d(\pi_{(P_0,Q_0)})_1$ is

$$\{Z \in (\mathcal{A}_0)_{ah} : ZP_0 = P_0Z \text{ and } ZQ_0 = Q_0Z\}.$$

This is the Banach-Lie algebra of $\pi_{(P_0,Q_0)}^{-1}(P_0,Q_0)$ (usually called the isotropy subgroup of the action at (P_0, Q_0)). It was described in Remark 4.3 using Halmos representation. It is clear, then, that the Banach Lie algebra of the isotropy group consists of matrices

$$\begin{pmatrix} X & 0 \\ 0 & X \end{pmatrix}$$

with $X^* = -X$ and $X\Gamma = \Gamma X$.

Let us prove that this space is complemented in the Banach-Lie algebra $(\mathcal{A}_0)_{ah}$, which (in this representation) consists of matrices of the form

$$\begin{pmatrix} X & Y \\ Y & Z \end{pmatrix}$$

where X, Y, Z commute with Γ, and satisfy the equation

$$C(X - Z) + 2SY = 0.$$

As remarked above, $E_{(P_0,Q_0)}$ is a conditional expectation between the von Neumann algebras $\mathcal{A}_0$ and the isotropy subalgebra $N(d(\pi_{(P_0,Q_0)})_1)$ at (P_0, Q_0). Therefore, the anti-Hermitian part of the nullspace $N(E_{(P_0,Q_0)}) \cap (\mathcal{A}_0)_{ah}$ is a supplement for the isotropy subalgebra. $\qquad\square$

In the general case, i.e., if $R(A^2 - 1)$ is not necessarily closed, we shall use the transitive action of $\mathcal{U}_{\mathcal{A}_0}$ to induce a differentiable structure in $\mathcal{D}_{A_0}$. Using Halmos representation, we know the explicit form of the isotropy subgroups of the action. In order to prove that $\mathcal{D}_{A_0}$ has a C^∞ structure, and that the maps $\pi_{(P_0,Q_0)}$ are submersions, we shall use a general result on quotients of unitary groups (see, for instance, [10]). In this result, it is required that H is a *Banach-Lie subgroup* of G in the following specific sense:

Definition 6.7. ([10] Definition 4.1) *Let G be a Banach-Lie group and H a subgroup of G. We say that H is a Banach-Lie subgroup of G if the following conditions are satisfied.*

1. *The subgroup H is endowed with a structure of Banach-Lie group whose underlying topology is the same as the relative topology of H in G.*

2. *The inclusion map $H \hookrightarrow G$ is smooth and the induced mapp between the Banach-Lie algebras $L(H) \to L(G)$ is an injective operator with closed range.*

3. *There exists a closed linear subspace M of $L(G)$ such that $L(H) \oplus M = L(G)$.*

Proposition 6.8. ([10] Theorem 4.19) *Let G be a Banach-Lie group, H a Banach-Lie subgroup of G and $\pi : G \to G/H$ the natural projection. Endow G/H with the quotient topology and consider the natural transitive action*

$$G \times G/H \to G/H, (g, kH) \mapsto gkH.$$

Then G/H has a structure of C^∞ manifold and the following conditions are satisfied:

1. *The mapping π is C^∞ and has C^∞ local cross sections near every point of G/H.*

2. *For every $g \in G$ the mapping*

$$G/H \to G/H \ , \quad kH \mapsto gkH$$

   is C^∞.

The isotropy subgroup $\mathcal{I}_{(P_0,Q_0)} := \{W \in \mathcal{U}_{\mathcal{A}_0} : (WP_0W^*, WQ_0W^*) = (P_0, Q_0)\} \subset \mathcal{U}_{\mathcal{A}_0}$ clearly satisfies the conditions of Definition 6.7 (the space $M := \mathbb{H}_{(P_0,Q_0)}$ given in (4), satisfies condition 3.).

Corollary 6.9. *If $A \in \mathcal{D}$, the space $\mathcal{D}_{A_0}$ inherits a C^∞ manifold structure from the quotient $\mathcal{U}_{\mathcal{A}_0}/\mathcal{I}_{(P_0,Q_0)}$, which makes $\pi_{(P_0,Q_0)}$ a C^∞ submersion.*

Remark 6.10. The topology that the quotient $\mathcal{U}_{\mathcal{A}_0}/\mathcal{I}_{(P_0,Q_0)}$ induces in $\mathcal{D}_{A_0}$ might be different from the ambient topology induced by $\mathcal{B}(\mathcal{H}_0) \times \mathcal{B}(\mathcal{H}_0)$. In other words, the identification is not necessarily a homeomorphism between these topologies.

7 A reductive structure for $\mathcal{D}_{A_0}$

Recall the supplement $\mathbb{H}_{(P_0,Q_0)}$ of $L(\mathcal{I}_{(P_0,Q_0)}) = N(E_{(P_0,Q_0)}) \cap (\mathcal{A}_0)_{ah}$ in $(\mathcal{A}_0)_{ah}$, defined in 6.3,

$$\mathbb{H}_{(P_0,Q_0)} = \{ \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix} : Y^* = -Y, Y\tau \text{ bounded} \}.$$

This distribution of subspaces $\mathcal{D}_{A_0} \ni (P_0, Q_0) \mapsto \mathbb{H}_{(P_0,Q_0)}$ is what in differential geometry is called a *reductive structure* for $\mathcal{D}_{A_0}$, meaning that it satisfies the following conditions:

- The subspace $\mathbb{H}_{(P_0,Q_0)}$ is invariant under the inner action of $\mathcal{I}_{(P_0,Q_0)}$: if $W \in \mathcal{I}_{(P_0,Q_0)}$ and $Z \in \mathbb{H}_{(P_0,Q_0)}$ then $W \cdot Z \in \mathbb{H}_{(P_0,Q_0)}$.

- The distribution of supplements $\mathcal{D}_{A_0} \ni (P', Q') \mapsto \mathbb{H}_{(P',Q')}$ is C^∞. This means that if $\mathbf{P}_{\mathbb{H}_{(P',Q')}}$ denotes the idempotent (real) linear map acting in $(\mathcal{A}_0)_{ah}$ corresponding to the projection to the first component in the decomposition $(\mathcal{A}_0)_{ah} = \mathbb{H}_{(P',Q')} \oplus L(\mathcal{I}_{(P',Q')})$, then the map

$$\mathcal{D}_{A_0} \ni (P', Q') \mapsto \mathbf{P}_{\mathbb{H}_{(P',Q')}} \in \mathcal{B}((\mathcal{A}_0)_{ah})$$

  is C^∞.

In our case, the fact that the supplement is the (anti-Hermitian part) of the nullspace of an $L(\mathcal{I}_{(P_0,Q_0)})$-valued conditional expectation, implies the first property. Also note that $\mathbf{P}_{\mathbb{H}_{(P',Q')}} = \left(Id - E_{(P_0,Q_0)} \right)|_{\mathcal{B}((\mathcal{A}_0)_{ah})}$.

Each pair (P', Q') gives rise to a conditional expectation, which due to Remark 6.2, depends only on the pair. We must show that if the pairs (P', Q') vary smoothly, then so do the maps $E_{(P',Q')}$. The map $B \mapsto V$, via the polar decomposition $B = V|B|$, in general is not continuous, much less smooth. However, in the unitary orbit of (P_0, Q_0), it is smooth. Indeed, note that, locally (for (P', Q') close to (P_0, Q_0)), the unitary U in $\mathcal{A}_0$ such that $(U P_0 U^*, U Q_0 U^*) = (P', Q')$ can be chosen as a smooth map in the arguments P', Q', by means of smooth local cross section for the submersion $\pi_{(P_0,Q_0)}$. Then, if we denote (as in Remark 6.2), $B = Q_0 - P_0 Q_0 P_0 + P_0^{\perp} Q_0 P_0^{\perp}$, and, accordingly, $B' = Q' - P'Q'P' - P'^{\perp}Q'P'^{\perp}$, and V and V' are the isometric parts in the corresponding polar decompositions of B and B', then

$$ UVU^* = V'. $$

Therefore, $E_{(P',Q')} = U E_{(P_0,Q_0)}(U^* \cdot U)U^*$.

Remark 7.1. Elements $Z \in \mathbb{H}_{(P_0,Q_0)}$ have symmetric spectrum, with symmetric multiplicity. Indeed, consider

$$ J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}. $$

Then $J^* = -J$, $J^2 = -1$ (a fortiori, J is a unitary in $\mathcal{H}_0$), and for any $Z \in \mathbb{H}_{(P_0,Q_0)}$,

$$ J \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix} = - \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix} J, $$

i.e., $JZJ^* = -ZJJ^* = -Z$.

Therefore, if one considers the symmetry $V = iJ$, it turns out that elements in $\mathbb{H}_{(P_0,Q_0)}$ are of the form $Z = r(E - F)$, for $r = \|Z\|$ and E, F orthogonal projections.

Remark 7.2. A reductive structure on a homogeneous space induces a linear connection (see [24], or [26] for an infinite dimensional setting). For instance, the geodesics can be explicitly computed (in terms of YC^{-1}). Indeed, note that, if $Z = \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix}$, then

$$ Z = Z_0 + Z_1, \quad \text{where } Z_0 = \begin{pmatrix} -Y\tau & 0 \\ 0 & Y\tau \end{pmatrix} \text{ and } Z_1 = \begin{pmatrix} 0 & Y \\ Y & 0 \end{pmatrix}, $$

and Z_0, Z_1 anti-commute. Thus,

$$ Z^2 = Z_0^2 + Z_1^2 = \begin{pmatrix} Y^2\tau^2 + Y^2 & 0 \\ 0 & Y^2\tau^2 + Y^2 \end{pmatrix} = (YC^{-1})^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, $$

and $Z^{2n} = (YC^{-1})^{2n} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$. Also

$$ Z^{2n+1} = Z^{2n} Z = \begin{pmatrix} -(YC^{-1})^{2n}Y\tau & (YC^{-1})^{2n}Y \\ (YC^{-1})^{2n}Y & (YC^{-1})^{2n}Y\tau \end{pmatrix} = \begin{pmatrix} -(YC^{-1})^{2n+1}S & (YC^{-1})^{2n+1}YC \\ (YC^{-1})^{2n+1}YC & (YC^{-1})^{2n+1}YS \end{pmatrix} $$

$$= (YC^{-1})^{2n+1} \begin{pmatrix} -S & C \\ C & S \end{pmatrix}.$$

Then,
$$e^{tZ} = \cosh(tYC^{-1}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \sinh(tYC^{-1}) \begin{pmatrix} -S & C \\ C & S \end{pmatrix}. \tag{6}$$

The operator $\Sigma = \begin{pmatrix} -S & C \\ C & S \end{pmatrix}$ is a symmetry: $\Sigma^* = \Sigma = \Sigma^{-1}$.

Note that, since YC^{-1} is anti-Hermitian, and cosh and sinh are, respectively, even and odd functions, then the first term of e^{tZ} is selfadjoint and the second is anti-Hermitian. Then,
$$e^{-tZ} = \cosh(tYC^{-1}) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \sinh(tYC^{-1}) \begin{pmatrix} -S & C \\ C & S \end{pmatrix}.$$

Therefore, the geodesic $\delta(t) = (e^{tZ}P_0 e^{-tZ}, e^{tZ}Q_0 e^{-tZ})$ can be explicitly computed.

Alternatively, let $D^* = D$ such that $YC^{-1} = iD$. Then, $\cosh(tYC^{-1}) = \cos(tD)$ and $\sinh(tYC^{-1}) = i\sin(tD)$.

Remark 7.3. The geodesics can be described in an intrinsic way, without reference to the Halmos frame of reference. To this effect, note that the matrix $\begin{pmatrix} -S & C \\ C & S \end{pmatrix}$ above is precisely $-J_0$, $(J_0 = sgn(A_0))$. Indeed:
$$A_0 = \begin{pmatrix} S^2 & -CS \\ -CS & -S^2 \end{pmatrix}, \quad A_0^2 = \begin{pmatrix} S^2 & 0 \\ 0 & S^2 \end{pmatrix} \quad \text{and} \quad |A_0| = \begin{pmatrix} S & 0 \\ 0 & S \end{pmatrix}.$$

Next, note that
$$P_0 Z P_0 J_0 P_0 = \begin{pmatrix} -Y\tau & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} -S & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} YC^{-1} & 0 \\ 0 & 0 \end{pmatrix}.$$

Then
$$e^{tZ} = \cosh(tP_0 Z P_0 J_0 P_0) + \sinh(tP_0 Z P_0 J_0)J_0, \tag{7}$$
and similarly for e^{-tZ}.

With these expressions above, the bijectivity radius of the exponential map can be computed.

Theorem 7.4. *Let $(P_0, Q_0) \in \mathcal{D}_{A_0}$. Then the exponential map*

$$exp_{(P_0, Q_0)} : \{Z \in \mathbb{H}_{(P_0, Q_0)} : \|Z\| < \pi/2\} \to exp_{(P_0, Q_0)}\left(\{Z \in \mathbb{H}_{(P_0, Q_0)} : \|Z\| < \pi/2\}\right)$$

is a bijection, whose image

$$exp_{(P_0, Q_0)}\left(\{Z \in \mathbb{H}_{(P_0, Q_0)} : \|Z\| < \pi/2\}\right) \subset \mathcal{D}_{A_0}$$

is an open dense subset.

Proof. Let $v_1, v_2 \in (T\mathcal{D}_{A_0})_{(P_0,Q_0)}$ with the same exponential, i.e., if $Z_1, Z_2 \in \mathbb{H}_{(P_0,Q_0)}$ are the corresponding horizontal elements, $e^{Z_1} \cdot (P_0, Q_0) = e^{Z_2} \cdot (P_0, Q_0)$. Then, following the notations of the preceeding remark ($Z_j = \begin{pmatrix} -Y_j\tau & Y_j \\ Y_j & Y_j\tau \end{pmatrix}$ and $D_j = -iY_j\tau$, for $j = 1, 2$), there exists a unitary operator in $\mathcal{L}$ such that

$$e^{-Z_2}e^{Z_1} = \begin{pmatrix} W & 0 \\ 0 & W \end{pmatrix},$$

i.e., $\sin(D_2)W = \sin(D_1)$ and $\cos(D_2)W = \cos(D_1)$. If we suppose that $|v_j|_{(P_0,Q_0)} = \|D_j\| < \pi/2$, then the cosines are invertible, and, then, these identities imply that $D_1 = D_2$.

Let us prove that $exp_{(P_0,Q_0)}\left(\{Z \in \mathbb{H}_{(P_0,Q_0)} : \|Z\| < \pi/2\}\right)$ is an open dense subset of $\mathcal{D}_{A_0}$. In order to prove this fact, we shall use the alternative characterization for $\mathcal{D}_{A_0}$ given in Theorem 2.2. Namely, representing the elements of $\mathcal{D}_{A_0}$ as symmetries V which anti-commute with A_0 . Recall that the correspondence is given by $(P, Q) \leftrightarrow V$, where V is the symmetry (the isometric part) in the polar decomposition of $P + Q - 1$. Let us compute V_0, the symmetry corresponding to the pair (P_0, Q_0), using Halmos decomposition based on the pair (P_0, Q_0). Note that

$$P_0 + Q_0 - 1 = \begin{pmatrix} C^2 & CS \\ CS & -C^2 \end{pmatrix} , \quad (P_0 + Q_0 - 1)^2 = \begin{pmatrix} C^2 & 0 \\ 0 & C^2 \end{pmatrix},$$

so that

$$|P_0 + Q_0 - 1| = \begin{pmatrix} C & 0 \\ 0 & C \end{pmatrix} \quad \text{and} \quad V_0 = \begin{pmatrix} C & S \\ S & -C \end{pmatrix}.$$

Pick $Z \in \mathbb{H}_{(P_0,Q_0)}$, $Z = \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix}$. Recall that Y and $Y\tau$ commute with C and S, so we get

$$ZV_0 = \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix} \begin{pmatrix} C & S \\ S & -C \end{pmatrix} = \begin{pmatrix} 0 & -Y\tau S - YC \\ YC + Y\tau S & 0 \end{pmatrix}$$

and

$$V_0 Z = \begin{pmatrix} 0 & Y\tau S + YC \\ -YC - Y\tau S & 0 \end{pmatrix}.$$

That is, $ZV_0 = -V_0 Z$. Conversely, a similar computation shows that anti-Hermitian elements Z which anti-commute with V_0 belong to $\mathbb{H}_{(P_0,Q_0)}$. Pick a symmetry V which anti-commutes with A_0, such that $\|V - V_0\| < 2$. In general, any pair of symmetries lie at distance less or equal than 2. Therefore, the set of such V form an open dense subset of $\mathcal{D}_{A_0}$. We claim that these elements V belong to the image of the exponential $exp_{(P_0,Q_0)}$ restricted to $\{Z \in \mathbb{H}_{(P_0,Q_0)} : \|Z\| < \pi/2\}$. In [28], H. Porta and L.Recht proved that a symmetry V such that $\|V - V_0\| < 2$ is of the form $V = e^Z V_0 e^{-Z}$, where Z is an anti-Hermitian operator which anti-commutes with V_0, i.e., $Z \in \mathbb{H}_{P_0,Q_0}$, with $\|Z\| < \pi/2$. Also note that, since Z anti-commutes with V_0, $V_0 e^{-Z} = e^Z V_0$, thus,

$$V = e^{2Z} V_0 \quad \text{and} \quad e^{2Z} = VV_0.$$

Since $\|2Z\| < \pi$, Z can be obtained as $Z = \frac{1}{2}\log(VV_0)$ (log the unique logarithm for unitary operators U such that $\|U - 1\| < 2$; notice that $\|VV_0 - 1\| = \|V - V_0\| < 2$). Since V and V_0 anti-commute with A_0, VV_0 and its logarithm Z belong to $\mathcal{A}_0$.

The correspondence between symmetries and pairs in $\mathcal{D}_{A_0}$ is clearly equivariant: $UV_0 U^* \leftrightarrow U \cdot (P_0, Q_0)$, for $U \in \mathcal{U}_{A_0}$. Therefore, $(P, Q) = exp_{(P_0,Q_0)}(Z)$, where $Z \in \{Z \in \mathbb{H}_{(P_0,Q_0)} : \|Z\| < \pi/2\}$. $\qquad\square$

A similar argument (representing $\mathcal{D}_{A_0}$ as Davis' symmetries) allows one to prove that $exp_{(P_0,Q_0)}$ is globally onto. The proof is also a refinement of the argument which showed that the action of $\mathcal{U}_{A_0}$ on $\mathcal{D}_{A_0}$ is transitive. In fact, the result below implies the former.

Theorem 7.5. *Let $(P_0, Q_0) \in \mathcal{D}_{A_0}$. The exponential map*

$$exp_{(P_0,Q_0)} : \{Z \in \mathbb{H}_{(P_0,Q_0)} : \|Z\| \leq \pi/2\} \to \mathcal{D}_{A_0}$$

is onto.

Proof. Pick $(P, Q) \in \mathcal{D}_{A_0}$, and let $V = sgn(P + Q - 1)$ be its Davis' symmetry. Let V_0 be the symmetry corresponding to (P_0, Q_0). Denote by

$$\mathcal{H}_{+,+} = \{\xi \in \mathcal{H} : V_0\xi = \xi \text{ and } V\xi = \xi\}, \quad \mathcal{H}_{-,-} = \{\xi \in \mathcal{H} : V_0\xi = -\xi \text{ and } V\xi = -\xi\},$$

$$\mathcal{H}_{+,-} = \{\xi \in \mathcal{H} : V_0\xi = \xi \text{ and } V\xi = -\xi\} \text{ and } \mathcal{H}_{-,+} = \{\xi \in \mathcal{H} : V_0\xi = -\xi \text{ and } V\xi = \xi\}.$$

Recall the symmetry J_0 obtained as the sign of A_0, which anti-commutes with V and V_0. Note that, if $V_0\xi = \xi$, then $J_0\xi \neq 0$ satisfies $V_0 J_0 \xi = -J_0 V_0 \xi = -J_0 \xi$, and similarly for ξ such that $V_0\xi = -\xi$. The analogous property holds for V. Therefore, J_0 maps $\mathcal{H}_{+,+}$ into $\mathcal{H}_{-,-}$ and viceversa, and it is one to one between these subspaces. The same happens for $\mathcal{H}_{+,-}$ and $\mathcal{H}_{-,+}$. In [2], it was proven that there exists a geodesic of the Grassmann manifold of $\mathcal{H}$ (=space of symmetries of $\mathcal{H}$) joining V and V_0 if and only if

$$\dim(\mathcal{H}_{+,-}) = \dim(\mathcal{H}_{-,+}).$$

It follows that there exists Z, $Z^* = -Z$, $\|Z\| \leq \pi/2$, which is co-diagonal with respecto to V_0, such that $e^Z V_0 e^{-Z} = V$. Such Z may not be unique (if the dimensions above are non zero). To finish the proof, we must show that we can find one of these Z in $\mathcal{A}_0$. In [2] it was shown that all such Z are reduced by the decomposition

$$\mathcal{H}_0 = \mathcal{H}_0' \oplus \mathcal{H}_0'' \oplus \mathcal{H}_0^0,$$

where $\mathcal{H}_0' = \mathcal{H}_{+,+} \oplus \mathcal{H}_{-,-}$, $\mathcal{H}_0' = \mathcal{H}_{+,-} \oplus \mathcal{H}_{-,+}$ and $\mathcal{H}_0^0$ is the orthogonal complement to the sum of these two. Both J_0 and A_0 are reduced by this decomposition. Clearly, also V and V_0 are also reduced.

Z is trivial in $\mathcal{H}_0'$ ([2]). In $\mathcal{H}_0''$, we choose $Z|_{\mathcal{H}_0''} = i\pi/2 J_0|_{h_0''}$. The exponential of this anti-Hermitian operator yields a unitary operator intertwinig $V_0|_{\mathcal{H}_0''}$ and $V|_{\mathcal{H}_0''}$. In the third subspace, which, in fact, is the generic part of the pair V_0, V, the exponent $Z|_{\mathcal{H}_0^0}$ is obtained uniquely as the logarithm (see [2])

$$Z|_{\mathcal{H}_0^0} = \log(\mathbf{S}V_0|_{\mathcal{H}_0^0}),$$

where $\mathbf{S}$ is the Davis symmetry of the reductions $V_0|_{\mathcal{H}_0^0}$ and $V|_{\mathcal{H}_0^0}$ to their generic part:

$$\mathbf{S} = sgn(-\frac{1}{2}\{V_0 + V\}|_{\mathcal{H}_0^0}).$$

$\{V + V_0\}|_{\mathcal{H}_0^0}$ anti-commutes with $A_0|_{\mathcal{H}_0^0}$, then so does its unitary part $\mathbf{S}$ in the polar decomposition. It follows that $\mathbf{S}V_0$ commutes with $A_0|_{\mathcal{H}_0^0}$. Therefore, our choice of Z belongs to $\mathcal{A}_0$ and the proof is complete . $\qquad\square$

8 A Hopf-Rinow theorem for $\mathcal{D}_{A_0}$

In [17], Durán, Mata-Lorenzo and Recht introduced a Finsler metric for homogeneous spaces $\mathcal{U}_\mathcal{A}/\mathcal{U}_\mathcal{B}$ which are obtained as the quotient of the unitary group of a C*-algebra by the unitary group of a C*-subalgebra $\mathcal{B} \subset \mathcal{A}$. Denote by

$$\pi : \mathcal{U}_\mathcal{A} \to \mathcal{U}_\mathcal{A}/\mathcal{U}_\mathcal{B}$$

the quotient map. A tangent vector v at $[1]$ (the class of 1 in $\mathcal{U}_\mathcal{A}/\mathcal{U}_\mathcal{B}$) identifies with an element in the quotient of (real) Banach spaces $\mathcal{A}_{ah}/\mathcal{B}_{ah}$ ($\mathcal{A}_{ah}$ and $\mathcal{B}_{ah}$ denote the spaces of anti-Hermitian elements of $\mathcal{A}$ and $\mathcal{B}$, respectively). Thus $v = d(\pi)_1(z)$, for some $z \in \mathcal{A}_{ah}$; we shall say that z is a *lifting* of v. The Finsler norm $|v|_{[1]}$ is defined as the infimum

$$|v| = |v|_{[1]} = \inf\{\|z\| : z \in \mathcal{A}_{ah} \text{ is a lifting of } v\}.$$

Or equivalently, if z_0 is an arbitrary lifting of v,

$$|v| = \inf\{\|z_0 + y\| : y \in \mathcal{B}_{ah}\}.$$

The metric is carried over the entire tangent bundle by means of the left action of $\mathcal{U}_\mathcal{A}$ on $\mathcal{U}_\mathcal{A}/\mathcal{U}_\mathcal{B}$. Thus, the metric so defined is invariant under this action. A lifting z_0 which achieves this infimum is called a *minimal lifting*. Minimal liftings may not exist (see for instance [11] for a nice example), and may fail to be unique (see [8] for a finite dimensional example), but if $\mathcal{B} \subset \mathcal{A}$ are von Neumann algebras, they do exist [18]. We shall apply this theory to $\mathcal{D}_{A_0}$, where the algebras are indeed von Neumann algebras (being defined as commutant of selfadjoint operators). The main result in [17] states that if $z_0 \in \mathcal{A}_{ah}$ is a minimal lifting of v, then

$$\delta(t) = [e^{tz}]$$

is a curve of minimal length for time $|t| \leq \frac{\pi}{2|v|}$.

In our context $L(\mathcal{I}_{(P_0,Q_0)})$ is the space of anti-Hermitian elements of the von Neumann algebra of operators which commute with P_0 and Q_0, inside the von Neumann algebra of operators which commute with $A_0 = P_0 - Q_0$. The fact that $\mathbb{H}_{(P_0,Q_0)}$ is a supplement for $L(\mathcal{I}_{(P_0,Q_0)})$ in $(\mathcal{A}_0)_{ah}$, implies that any tangent vector v in $T(\mathcal{D}_{A_0})_{(P_0,Q_0)}$ has a lifting $Z \in \mathbb{H}_{(P_0,Q_0)}$. We claim that this lifting is minimal. To prove this we need the following result:

Proposition 8.1. (Proposition 5.2 of [17], see also Theorem 2.2 of [8]) *Let $\mathcal{B} \subset \mathcal{A}$ be C*-algebras, and let $Z \in \mathcal{A}_{ah}$. Suppose that there exists a state ψ in $\mathcal{A}$ with the following properties:*

1. $\psi(Z^2) = -\|Z\|^2$.

2. For any $Y \in \mathcal{B}$, $\psi(YZ) = 0$.

Then Z is a minimal lifting (i.e., $\|Z\| \leq \|Z + D\|$, for all $D \in \mathcal{B}_{ah}$).

Lemma 8.2. *Let $Z^* = -Z$, with matrix*

$$Z = \begin{pmatrix} X & Y \\ Y & -X \end{pmatrix}$$

such that X and Y commute. Then

$$\|Z\| \leq \|Z + D\|,$$

for any $D^* = -D$ of the form $D = \begin{pmatrix} D' & 0 \\ 0 & D' \end{pmatrix}$.

Proof. We use Proposition 8.1 above. Note that

$$Z^2 = \begin{pmatrix} X^2 + Y^2 & XY - YX \\ YX - XY & X^2 + Y^2 \end{pmatrix} = \begin{pmatrix} X^2 + Y^2 & 0 \\ 0 & X^2 + Y^2 \end{pmatrix}.$$

Note, also, that $X^2 + Y^2 \leq 0$, and that $\|Z^2\| = \|X^2 + Y^2\|$. There exists a state ψ_0 in $\mathcal{B}(\mathcal{L})$ such that $\psi_0(X^2 + Y^2) = -\|X^2 + Y^2\|$. Let τ the positive unital linear map

$$\tau(\begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix}) = \frac{1}{2}\{T_{11} + T_{22}\}.$$

Then $\psi = \psi_0 \circ \tau$ is a state in $\mathcal{B}(\mathcal{H})$ such that

1.
$$\psi(Z^2) = \psi_0(X^2 + Y^2) = -\|X^2 + Y^2\|.$$

2. For $D^* = -D$ of the form, $D = \begin{pmatrix} D' & 0 \\ 0 & D' \end{pmatrix}$,

$$\tau(ZB) = \tau(\begin{pmatrix} XD' & YD' \\ YD' & -XD' \end{pmatrix}) = 0.$$

Then $\psi(ZB) = 0$.

$\square$

If $v \in T(\mathcal{D}_{A_0})_{(P_0,Q_0)}$, we shall denote by $Z_v \in \mathbb{H}_{(P_0,Q_0)}$ the unique lifting of v which belongs to $\mathbb{H}_{(P_0,Q_0)}$. Note that the mapping $v \mapsto Z_v$ is a linear isomorphism (it is the inverse of $d(\pi_{(P_0,Q_0)})_1$ restricted to $\mathbb{H}_{(P_0,Q_0)}$).

The following result follows:

Theorem 8.3. *Let* $(P_0, Q_0) \in \mathcal{D}_{A_0}$ *and* $v \in T(\mathcal{D}_{A_0})_{(P_0,Q_0)}$ *such that* $Z_v = \begin{pmatrix} Y\tau & Y \\ Y & -Y\tau \end{pmatrix}$. *Then*

$$\delta(t) = e^{tZ_v} \cdot (P_0, Q_0) = (e^{tZ_v} P_0 e^{-tZ_v}, e^{tZ_v} Q_0 e^{-tZ_v})$$

is a minimal geodesic for the Finsler metric defined above, up to time

$$|t| \leq \frac{\pi}{2|v|_{(P_0,Q_0)}} = \frac{\pi}{2\|Z_v\|} = \frac{\pi}{2\|Y^2 + (Y\tau)^2\|^{1/2}}.$$

Proof. Recall from Remark 6.4, that for any $Z_v \in \mathbb{H}_{(P_0,Q_0)}$, all entries in the matrix commute. Thus, Lemma 8.2 holds and Z_v has minimal norm among all perturbations with elements in the isotropy algebra. $\square$

The fact that the minimal liftings belong to the linear space $\mathbb{H}_{(P_0,Q_0)}$ implies, additionally, the following result, which can be regarded as a Hopf-Rinow Theorem for the space $\mathcal{D}_{A_0}$ (in a context which is far from being Riemannian). Note that the pairs in $\mathcal{D}_{A_0}$ which can be reached from any given pair (P_0, Q_0) by a *unique* minimal geodesic, is an open dense subset of $\mathcal{D}_{A_0}$. Theorem 7.4 and Theorem 8.3 imply

Corollary 8.4. *Let* $(P, Q), (P_0, Q_0) \in \mathcal{D}_{A_0}$, *and let* V *and* V_0 *be the corresponding Davis' symmetries:* $V = sgn(P + Q - 1)$, $V_0 = sgn(P_0 + Q_0 - 1)$. *If* $\|V - V_0\| < 2$, *then there exists a unique element* $Z = Z_{(P_0, Q_0)}(P, Q)$, *which is a* C^∞ *map in terms of the arguments* $(P_0, Q_0), (P, Q)$, *such that*

$$(P, Q) = e^Z \cdot (P_0, Q_0).$$

The geodesic $\delta(t) = e^{tZ} \cdot (P_0, Q_0)$ *has minimal length among all piecewise smooth curves in* $\mathcal{D}_{A_0}$ *joining* (P_0, Q_0) *with* (P, Q).

Proof. From the proof of of Theorem 7.4, it is clear that $\|Z\| < \pi/2$. Then, by Theorem 8.3, δ is minimal up to time $t = 1$, and $\delta(1) = (P, Q)$. $\qquad\square$

Using now Theorem 7.5, one can show that (dropping the uniqueness condition), any pair of elements in $\mathcal{D}_{A_0}$ can be joined by a minimal geodesic.

Corollary 8.5. *Let* $(P_0, Q_0), (P, Q) \in \mathcal{D}_{A_0}$. *Then there exists a minimal geodesic of* $\mathcal{D}_{A_0}$, *of length less or equal than* $\pi/2$, *which joins them.*

Proof. The existence of a geodesic δ, parametrized in the interval $[0, 1]$, joining (P_0, Q_0) and (P, Q) is a direct consequence of Theorem 7.5. Also, it is clear there that the norm of the exponent Z (=the length of the geodesic) is less or equal than $\pi/2$. Let us prove, by a standard geometric argument, that it must be minimal. Clearly, by the above result, we must consider the case when the length $\ell(\delta)$ of δ is $\pi/2$. Suppose there is a curve $\gamma \in \mathcal{D}_{A_0}$ joining (P_0, Q_0) of length $\pi/2 - r$, for some $r > 0$. Pick t_0 such that $(1 - t_0)\pi/2 = \ell(\delta|_{[t_0, 1]}) < r/2$. Then the length of the curve obtained by adjoining to γ the curve $\delta|_{[t_0, 1]}$ reversed, gives a curve joining (P_0, Q_0) and $\delta(t_0)$, of length strictly less than $t_0\pi/2$. This is a contradiction, since by the above Corollary, the curve $\delta|_{[0, t_0]}$ has minimal length $t_0\pi/2$. $\qquad\square$

Remark 8.6. Note that if $Z = \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix}$ is a horizontal element, then also YC^{-1} is bounded. Recall that X, Y, Z commute with C, S, τ, the equality $C(Z - X) = 2SY$ implies that

$$C^2(Z - X)^2 4S^2 Y^2 = 4Y^2 - 4C^2 Y^2, \quad \text{i.e.,} \quad C^2 A = Y^2,$$

where $A = \frac{1}{4}((X - Z)^2 + Y^2)$ is bounded. It follows that $Y^2 C^{-2}$ is bounded, thus its square root $|Y|c^{-1}$ is bounded, and then also YC^{-1} is bounded. Moreover, if $Z = \begin{pmatrix} -Y\tau & Y \\ Y & Y\tau \end{pmatrix}$, then

$$\|Z\| = \|Y^2 \tau^2 + Y^2\|^{1/2} = \|YC^{-1}\|.$$

Therefore, if $\|Z\| \leq \pi/2$, the distance between (P_0, Q_0) and $e^Z \cdot (P_0, Q_0)$ equals $\|YC^{-1}\|$.

References

[1] Amrein, W.O.; Sinha, K.B., On pairs of projections in a Hilbert space, Linear Algebra Appl. 208/209 (1994), 425–435.

[2] Andruchow, E.; Operators which are the difference of two projections, J. Math. Anal. Appl. 420 (2014), no. 2, 1634-1653.

[3] Andruchow, E.; Chiumiento, E.; Larotonda, G.; Geometric significance of Toeplitz kernels, J. Funct. Anal. (to appear), preprint arXiv:1608.05737

[4] Andruchow, E.; Corach, G., Uncertainty principle and geometry of the infinite Grassmann manifold, Studia Math. (to appear), preprint arXiv:1701.03733.

[5] Andruchow, E.; Corach, G., Schmidt decomposable products of projections, Integral Equations Operator Theory 79 (2018), 79-100.

[6] Andruchow, E.; Corach, G., Curves of projections an operator inequalities, Proceedings of the 26th International Conference on Operator Theory, Timisoara June 17th-July 2nd, 2016 (in press).

[7] Andruchow, E.; Corach, G.; Stojanoff, D. Projective spaces of a C^*-algebra, Integral Equations Operator Theory 37 (2000), 143–168.

[8] Andruchow, E.; Mata-Lorenzo, L. E.; Recht, L.; Mendoza, A.; Varela, A., Infinitely many minimal curves joining arbitrarily close points in a homogeneous space of the unitary group of a C^*-algebra, Rev. Un. Mat. Argentina 46 (2005), no. 2, 113–120 (2006).

[9] Avron, J.; Seiler, R.; Simon, B., The index of a pair of projections. J. Funct. Anal. 120 (1994), 220–237.

[10] Beltita, D., Smooth homogeneous structures in operator theory. Chapman & Hall/CRC Monographs and Surveys in Pure and Applied Mathematics, 137. Chapman & Hall/CRC, Boca Raton, FL, 2006.

[11] Bottazzi, T.; Varela, A., Minimal length curves in unitary orbits of a Hermitian compact operator, Differential Geom. Appl. 45 (2016), 1–22

[12] Corach, G.; Porta, H.; Recht, L., The geometry of spaces of projections in C^*-algebras, Adv. Math. 101 (1993), no. 1, 59–77.

[13] Davis, C., Separation of two linear subspaces, Acta Sci. Math. Szeged 19 (1958) 172–187.

[14] Deutsch, F., Best approximation in inner product spaces, Springer-Verlag, New York, 2001.

[15] Dixmier, J., Position relative de deux variétés linéaires fermées dans un espace de Hilbert, Revue Sci. 86 (1948), 387-399.

[16] Douglas, R. G., Banach algebra techniques in operator theory. Second edition. Graduate Texts in Mathematics, 179. Springer-Verlag, New York, 1998.

[17] Durán, C. E.; Mata-Lorenzo, L. E.; Recht, L., Metric geometry in homogeneous spaces of the unitary group of a C^*-algebra. I. Minimal curves, Adv. Math. 184 (2004), 342–366

[18] Durán, C. E.; Mata-Lorenzo, L. E.; Recht, L., Metric geometry in homogeneous spaces of the unitary group of a C^*-algebra. II. Geodesics joining fixed endpoints. Integral Equations Operator Theory 53 (2005), 33–50.

[19] Folland, G. B.; Sitaram, A., The uncertainty principle: a mathematical survey, J. Fourier Anal. Appl. 3 (1997), 207–238.

[20] Friedrichs, K., On certain inequalities and characteristic value problems for analytic functions and for functions of two variables, Trans. Amer. Math. Soc. 41 (1937), 321–364.

[21] Halmos, P. R., Two subspaces, Trans. Amer. Math. Soc. 144 (1969), 381–389.

[22] Havin, V.; Jöricke, B., The uncertainty principle in harmonic analysis. Ergebnisse der Mathematik und ihrer Grenzgebiete (3) 28. Springer-Verlag, Berlin, 1994.

[23] Kato, T., Perturbation theory for linear operators. Reprint of the 1980 edition. Classics in Mathematics. Springer-Verlag, Berlin, 1995.

[24] Kobayashi, S.; Nomizu, K., Foundations of differential geometry. Vol. I. Reprint of the 1963 original. Wiley Classics Library. A Wiley-Interscience Publication. John Wiley & Sons, Inc., New York, 1996.

[25] Lenard, A., The numerical range of a pair of projections, J. Functional Analysis 10 (1972), 410–423.

[26] Mata-Lorenzo, L. E.; Recht, L., Infinite-dimensional homogeneous reductive spaces. Acta Cient. Venezolana 43 (1992), 76–90.

[27] Nikol'skii, N. K., Treatise on the shift operator. Spectral function theory. Grundlehren der Mathematischen Wissenschaften, 273. Springer-Verlag, Berlin, 1986.

[28] Porta, H.; Recht, L., Minimality of geodesics in Grassmann manifolds, Proc. Amer. Math. Soc. 100 (1987), 464–466.

[29] Raeburn, I., The relationship between a commutative Banach algebra and its maximal ideal space, J. Funct. Anal. 25 (1977), 366–390.

[30] Shi, W.; Ji, G.; Du, H., Pairs of orthogonal projections with a fixed difference, Linear Algebra Appl. 489 (2016), 288–297.

[31] Slepian, D.; Pollak, H. O., Prolate spheroidal wave functions, Fourier analysis and uncertainty. I, Bell System Tech. J. 40 (1961), 43–63.

Esteban Andruchow
Instituto de Ciencias, Universidad Nacional de Gral. Sarmiento,
J.M. Gutierrez 1150, (1613) Los Polvorines, Argentina
and Instituto Argentino de Matemática, 'Alberto P. Calderón', CONICET,
Saavedra 15 3er. piso, (1083) Buenos Aires, Argentina.
e-mail: eandruch@ungs.edu.ar

Gustavo Corach
Instituto Argentino de Matemática, 'Alberto P. Calderón', CONICET,
Saavedra 15 3er. piso, (1083) Buenos Aires, Argentina,

and Depto. de Matemática, Facultad de Ingeniería, Universidad de Buenos Aires, Argentina.
e-mail: gcorach@fi.uba.ar

Lázaro Recht
Departamento de Matemática P y A, Universidad Simón Bolívar
Apartado 89000, Caracas 1080A, Venezuela
e-mail: recht@usb.ve