

# The Integral Representation of the Elementary Solution of the $n$ -Dimensional Ultrahyperbolic Klein-Gordon Operator Iterated $k$ -Times.

SUSANA ELENA TRIONE

SUMMARY. The kernel  $G_\alpha(P \pm i0, m, n)$  (formula (I,1;3)) can be expressed by an integral representation by means of a simple (symbolic) Fourier integral (cf. Theorem 1, form. (II,1;7)). Theorem 2 expresses the useful particular case  $\alpha=2k$  where  $G_{2k}(P \pm i0, m, n)$  is the elementary (causal, anticausal) solution of the Klein-Gordon operator (cf. [1], p. 38, form. (II, 2;12)). The particular case of the Theorem 2, corresponding to  $n=4$ ,  $q=1$ ,  $k=1$  (cf. Theorem 3) is very interesting because the formula (III,1;4), when we adopt the plus sign, has been established by Bogoliubov-Parasiuk (cf. [5], p. 233) and plays an essential role in the theory, due to these authors, of causal distributional products. Moreover, we know that  $G_0(P \pm i0, m, n) = \delta$  (cf. [1], p. 37, form. (II,2;7)); therefore by taking limits in both members of (II,1;7) for  $k \rightarrow 0$ , we obtain an integral representation of the  $\delta$  measure of Dirac. We know that  $G_\alpha(P \pm i0, m=0, n) = H_\alpha(P \pm i0, n)$  (cf. [1], p. 49, Th.1, form. (II,7;1)). We give in V.1 an integral representation of  $H_\alpha(P \pm i0, n)$  by means of an inverse Laplace transform (cf. Theorem 4, formula (V,1;9)). Finally, we study the particular cases  $H_{2k}(P \pm i0, n)$  and  $H_0(P \pm i0, n)$  (cf. (V,1;11) and Note).

## I.1. Introduction.

Let  $x = (x_1, x_2, \dots, x_n)$  be a point of the  $n$ -dimensional Euclidean space  $\mathbb{R}^n$ . Consider a non-degenerate quadratic form in  $n$  variables of the form

$$P = P(x) = x_1^2 + \dots + x_p^2 - x_{p+1}^2 - \dots - x_{p+q}^2, \quad (\text{I, 1; 1})$$

where  $n = p + q$ . The distribution  $(P \pm i0)^\lambda$  is defined by

$$(P \pm i0)^\lambda = \lim_{\varepsilon \rightarrow 0} \{P \pm i\varepsilon|x|^2\}^\lambda, \quad (\text{I}, 1; 2)$$

where  $\varepsilon = 0$ ,  $|x|^2 = x_1^2 + \dots + x_n^2$ ,  $\lambda \in \mathcal{C}$ .

The distributions  $(P \pm i0)^\lambda$  are an important contribution of Gelfand (cf. [2], p. 274).

The distributions  $(P \pm i0)^\lambda$  are analytic in  $\lambda$  everywhere except at  $\lambda = -\frac{n}{2} - k$ ,  $k = 0, 1, \dots$ , where they have simple poles (cf. [2], p. 275).

Let  $G_\alpha(P \pm i0, m, n)$  be the causal (anticausal) distribution defined by

$$G_\alpha(P \pm i0, m, n) = H_\alpha(m, n)(P \pm i0)^{\frac{1}{2}(\frac{\alpha-n}{2})} K_{\frac{n-\alpha}{2}} \left( \sqrt{m^2(P \pm i0)} \right), \quad (\text{I}, 1; 3)$$

where  $m$  is a positive real number,  $\alpha \in \mathcal{C}$ ,  $K_\mu$  designates the well-known modified Bessel function of the third kind (cf. [3], p. 78, formulae (6) and (7)):

$$K_\nu(z) = \frac{\pi}{2} \frac{I_{-\nu}(z) - I_\nu(z)}{\sin \pi \nu}, \quad (\text{I}, 1; 4)$$

$$I_\nu(z) = \sum_{m=0}^{\infty} \frac{\left(\frac{z}{2}\right)^{2m+\nu}}{m! \Gamma(m + \nu + 1)}, \quad (\text{I}, 1; 5)$$

$H_\alpha(m, n)$  is the constant defined by

$$H_\alpha(m, n) = \frac{e^{i\frac{\pi}{2}\alpha} 2^{1-\frac{\alpha}{2}} e^{\pm \frac{\pi}{2}qi} (m^2)^{\frac{1}{2}(\frac{n-\alpha}{2})}}{(2\pi)^{\frac{n}{2}} \Gamma\left(\frac{\alpha}{2}\right)}, \quad (\text{I}, 1; 6)$$

$\nu \neq \text{integer}$ .

## II.1. The integral representation of the causal (anticausal) distribution $G_\alpha(P \pm i0, m, n)$ .

We shall give a representation of the kernel  $G_\alpha(P \pm i0, m, n)$  (cf. (I,1;3)) by means of a simple (symbolic) Fourier integral.

We start from the formula (40), p. 283, [4], vol I,

$$\begin{aligned} \int_0^\infty e^{-pt} \frac{1}{2} t^{-\nu-1} e^{-\frac{\alpha}{t}} dt &\stackrel{\text{def}}{=} f(p, \alpha, \nu) \\ &= \alpha^{-\frac{\nu}{2}} p^{\frac{\nu}{2}} K_\nu \left[ 2\alpha^{\frac{1}{2}} p^{\frac{1}{2}} \right], \end{aligned} \quad (\text{II}, 1; 1)$$

valid for  $\text{Re } p > 0$  and  $\text{Re } \alpha > 0$

By making the change of parameters:

$$\begin{aligned} p &= \pm i (m^2 \mp i\varepsilon), \\ \alpha &= \mp \frac{i}{4} (P \pm i\varepsilon |x|^2), \end{aligned}$$

where  $\varepsilon > 0$ ,  $m > 0$  and  $P = P(x)$  defined by (I,1;1); formula (II,1;1) can be written

$$\begin{aligned} &f\left(\pm i (m^2 \mp i\varepsilon), \mp \frac{i}{4} (P \pm i\varepsilon |x|^2), \nu\right) \\ &= 2^\nu e^{\pm i \frac{\pi}{2} \nu} (m^2 \mp i\varepsilon)^{\frac{\nu}{2}} (P \pm i\varepsilon |x|^2)^{-\frac{\nu}{2}} \cdot K_\nu \left[ \sqrt{(m^2 \pm i\varepsilon) (P \pm i\varepsilon |x|^2)} \right]. \end{aligned} \quad (\text{II}, 1; 2)$$

We obtain, taking limits in both members of (II,1;2) for  $\varepsilon \rightarrow 0$ ,

$$\begin{aligned} &\lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{\pm i (m^2 \mp i\varepsilon) t} t^{-\nu-1} e^{\mp \frac{i}{4t} (P \pm i\varepsilon |x|^2)} dt \\ &= 2^{\nu+1} e^{\pm i \frac{\pi}{2} \nu} m^\nu (P \pm i0)^{-\frac{\nu}{2}} K_\nu \left[ \sqrt{m^2 (P \pm i0)} \right]. \end{aligned} \quad (\text{II}, 1; 3)$$

By putting  $\nu = \frac{n-\alpha}{2}$ , we have

$$\begin{aligned} &\lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{\pm i (m^2 \mp i\varepsilon) t} t^{-\left(\frac{n-\alpha}{2}\right)-1} e^{\mp \frac{i}{4t} (P \pm i\varepsilon |x|^2)} dt \\ &= 2^{\frac{n-\alpha}{2}+1} e^{\pm i \frac{\pi}{2} \left(\frac{n-\alpha}{2}\right)} m^{\frac{n-\alpha}{2}} (P \pm i0)^{-\frac{\left(\frac{n-\alpha}{2}\right)}{2}} \cdot K_{\frac{n-\alpha}{2}} \left( \sqrt{m^2 (P \pm i0)} \right). \end{aligned} \quad (\text{II}, 1; 4)$$

By multiplying both members of (II,1;4) by adequate constants we can write, equivalently,

$$\begin{aligned} &C_\pm(n, \alpha, q) \lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{\pm i (m^2 \mp i\varepsilon) t} t^{-\left(\frac{n-\alpha}{2}\right)-1} e^{\mp \frac{i}{4t} (P \pm i\varepsilon |x|^2)} dt \\ &= H_\alpha(m, n) P(\pm i0)^{\frac{1}{2} \left(\frac{\alpha-n}{2}\right)} K_{\frac{n-\alpha}{2}} \left( \sqrt{m^2 (P \pm i0)} \right), \end{aligned} \quad (\text{II}, 1; 5)$$

where the constants  $C_{\pm}(n, \alpha, q)$  are

$$C_{\pm}(n, \alpha, q) = \frac{2^{1-\frac{n}{2}} e^{i\frac{\pi}{2}\alpha} e^{\mp i\frac{\pi}{2}(\frac{n-\alpha}{2})} e^{\pm i\frac{\pi}{2}q}}{(2\pi)^{\frac{n}{2}} \Gamma(\frac{\alpha}{2})}. \quad (\text{II}, 1; 6)$$

From the definitory formula (I,1;3) the right-hand member of (II,1;5) is, precisely,  $G_{\alpha}(P \pm i0, m, n)$ . Therefore, we have prove the following theorem.

**Theorem 1.** *The causal (anticausal) distribution kernel  $G_{\alpha}(P \pm i0, n)$  admits the integral representation*

$$G_{\alpha}(P \pm i0, m, n) = C_{\pm}(n, \alpha, q) \lim_{\varepsilon \rightarrow 0} \int_0^{\infty} e^{\pm i(m^2 \mp i\varepsilon)t} t^{-(\frac{n-\alpha}{2})-1} e^{\mp \frac{i}{4t}(P \pm i\varepsilon|x|^2)} dt, \quad (\text{II}, 1; 7)$$

where the constants  $C_{\pm}(n, \alpha, q)$  are

$$C_{\pm}(n, \alpha, q) = \frac{2^{1-\frac{n}{2}} e^{i\frac{\pi}{2}\alpha} e^{\mp i\frac{\pi}{2}(\frac{n-\alpha}{2})} e^{\pm i\frac{\pi}{2}q}}{(2\pi)^{\frac{n}{2}} \Gamma(\frac{\alpha}{2})}. \quad (\text{II}, 1; 8)$$

### III.1. Particular cases.

The distributional functions  $G_{2k}(P \pm i0, m, n)$  where  $n = \text{integer} \geq 2$  and  $k = 1, 2, \dots$  (called “Feynman distributions”) are elementary causal (anticausal) solutions of the ultrahyperbolic Klein-Gordon operator, iterated  $k$ -times (cf. Theorem 3, p. 38, formula (II,2;12) of [1]):

$$K^k \{G_{2k}(P \pm i0, m, n)\} = \delta,$$

where  $K^k$  is the  $n$ -dimensional ultrahyperbolic Klein-Gordon operator iterated  $k$ -times:

$$K^k = \left\{ \frac{\partial^2}{\partial x_1^2} + \dots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \dots - \frac{\partial^2}{\partial x_{p+q}^2} - m^2 \right\}^k, \quad (\text{III}, 1; 1)$$

$m > 0$  and  $p + q = n$ .

It is useful to register the particular case of the Theorem 1, corresponding to  $\alpha = 2k$ .

**Theorem 2.** *The elementary solutions  $G_{2k}(P \pm i0, m, n)$  admit the integral representation*

$$G_{2k}(P \pm i0, m, n) = C_{\pm}(n, k, q) \cdot \lim_{\varepsilon \rightarrow 0} \int_0^{\infty} e^{\pm i(m^2 \mp i\varepsilon)t} t^{k - \frac{n}{2} - 1} \cdot e^{\mp \frac{i}{4t}(P \pm i\varepsilon|x|^2)} dt, \quad (\text{III}, 1; 2)$$

where the constants  $C_{\pm}(n, k, q)$  are

$$C_{\pm}(n, k, q) = \frac{2^{1 - \frac{n}{2}} e^{\pm i \frac{\pi}{2} q} e^{\mp i \frac{\pi}{2} (\frac{n}{2} - k)}}{(2\pi)^{\frac{n}{2}} \Gamma(k)}. \quad (\text{III}, 1; 3)$$

Moreover, it is useful to register the particular case of this theorem, corresponding to  $n = 4$ ,  $q = 1$ ,  $k = 1$ .

**Theorem 3.** *The solutions  $G_2(P \pm i0, m, 4)$  of the Klein-Gordon operator admit the representation*

$$G_2(P \pm i0, m, 4) = \lim_{\varepsilon \rightarrow 0} \frac{1}{16\pi^2} \int_0^{\infty} e^{\pm i(m^2 \mp i\varepsilon)t} \frac{1}{t^2} e^{\mp \frac{1}{4t}(x_1^2 + x_2^2 + x_3^2 - x_4^2 \pm i\varepsilon|x|^2)} dt. \quad (\text{III}, 1; 4)$$

This last formula, when we adopt the plus sign, has been established by Bogoliubov-Parasiuk (cf. [5], p. 233) and plays an essential role in the theory, due to these authors, of causal distributional products.

#### IV.1. Note.

We know (cf. [1], p. 37, formula (II,2;7)), that

$$G_0(P \pm i0, m, n) = \delta.$$

Therefore, the thesis of Theorem 2, permit us, by taking limits in both members of (II,1;6) for  $k \rightarrow 0$ , obtain an integral representation of the  $\delta$ -measure of Dirac,

by means of a simple (symbolic) Fourier integral.

### V.1. The integral representation of the causal (anticausal) distribution $H_\alpha(P \pm i0, n)$ .

We define the causal (anticausal) distribution  $H_\alpha(P \pm i0, n)$  as follows:

$$H_\alpha(P \pm i0, n) = \frac{e^{\frac{i\pi\alpha}{2}} e^{\pm i\frac{\pi}{2}q} \Gamma\left(\frac{n-\alpha}{2}\right)}{2^\alpha \pi^{\frac{n}{2}} \Gamma\left(\frac{\alpha}{2}\right)} (P \pm i0)^{\frac{\alpha-n}{2}}, \quad (\text{V}, 1; 1)$$

where  $\alpha \in \mathcal{C}$ ,  $P$  is defined by (I,1;1) and  $q$  is the number of negative terms of the quadratic form  $P$ .

We know that  $G_\alpha(P \pm i0, m = 0, n) = H_\alpha(P \pm i0, n)$ . (cf. [1], page 49, Theorem 1, formula (II,7;1)).

The distributional functions  $H_\alpha$  are the causal (anticausal) analogues of the elliptic kernel of M. Riesz ([6], pp. 16-21) and have analogue properties, that we use to obtain causal (anticausal) solutions of the  $n$ -dimensional ultrahyperbolic operator, iterated  $k$ -times ( $k$  integer  $\geq 1$ ),

$$L^k = \left\{ \frac{\partial^2}{\partial x_1^2} + \dots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \dots - \frac{\partial^2}{\partial x_{p+q}^2} \right\}^k. \quad (\text{V}, 1; 2)$$

We shall give a representation of the kernel  $H_\alpha(P \pm i0, n)$  (cf. (V,1;1)) by means of an inverse Laplace transform.

We start from the formula (31), p. 264, [4], vol.1,

$$\begin{aligned} \int_0^\infty e^{-pt} \exp(-\alpha e^{-t}) dt &\stackrel{\text{def}}{=} f(\alpha, p) \\ &= \alpha^{-p} \gamma(p, \alpha), \end{aligned} \quad (\text{V}, 1; 3)$$

here (cf. [4], p. 387)

$$\begin{aligned} \gamma(\alpha, x) &= \int_0^x e^{-t} t^{\alpha-1} dt \\ &= \alpha^{-1} x^\alpha {}_1F_1(\alpha; \alpha + 1; -x), \end{aligned} \quad (\text{V}, 1; 4)$$

is the incomplete gamma function and  ${}_1F_1(a; c; z)$  is the Kummer's confluent hypergeometric series (cf. [4], p. 373).

By making the change of parameters

$$\begin{aligned}\alpha &= (P \pm i\varepsilon|x|^2), \\ p &= -\left(\frac{\alpha - n}{2}\right),\end{aligned}\tag{V, 1; 5}$$

where  $\varepsilon > 0$  and  $P = P(x) = x_1^2 + \dots + x_p^2 - x_{p+1}^2 - \dots - x_{p+q}^2$ ,  $p + q = n$ ; formula (V,1;3) can be written

$$\begin{aligned}f\left[(P \pm i\varepsilon|x|^2), -\left(\frac{\alpha - n}{2}\right)\right] &= \int_0^\infty e^{-(\frac{\alpha-n}{2})t} \exp [-(P \pm i\varepsilon|x|^2)e^{-t}] dt \\ &= (P \pm i\varepsilon|x|^2)^{\frac{\alpha-n}{2}} \gamma\left(-\frac{\alpha - n}{2}, P \pm i\varepsilon|x|^2\right).\end{aligned}\tag{V, 1; 6}$$

We obtain, taking limits in both members of (V,1;6), for  $\varepsilon \rightarrow 0$ ,

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{-(\frac{\alpha-n}{2})t} \exp [-(P \pm i\varepsilon|x|^2)e^{-t}] dt = (P \pm i0)^{\frac{\alpha-n}{2}} \gamma\left(\frac{n-\alpha}{2}, P \pm i0\right).\tag{V, 1; 7}$$

Taking into account the definitory formula of  $H_\alpha(P \pm i0, n)$  (cf. (V,1;1)), by multiplying both members of (V,1;7) by adequate constants, we can write, equivalently,

$$\begin{aligned}&C_\pm(n, q, \alpha)\gamma^{-1}\left(\frac{n-\alpha}{2}, P \pm i0\right) \cdot \lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{-\frac{\alpha-n}{2}t} \exp [-(P \pm i\varepsilon|x|^2)e^{-t}] dt \\ &= \frac{e^{i\frac{\pi}{2}\alpha} e^{\pm i\frac{\pi}{2}q} \Gamma\left(\frac{n-\alpha}{2}\right)}{2^\alpha \pi^{\frac{n}{2}} \Gamma\left(\frac{\alpha}{2}\right)} (P \pm i0)^{\frac{\alpha-n}{2}}.\end{aligned}\tag{V, 1; 8}$$

From the definitory formula (V,1;1) the right-hand member of (V,1;8) is, precisely,  $H_\alpha(P \pm i0, n)$ . Therefore, we have proved the following theorem

**Theorem 4.** *The kernels  $H_\alpha(P \pm i0)$  admit the integral representation*

$$H_\alpha(P \pm i0, n) = C_\pm(n, q, \alpha)\gamma^{-1}\left(\frac{n-\alpha}{2}, P \pm i0, n\right) \lim_{\varepsilon \rightarrow 0} \int_0^\infty e^{-(P \pm i\varepsilon|x|^2)t} e^{(\frac{\alpha-n}{2})e^{-t}} dt,\tag{V, 1; 9}$$

where the constants  $C_{\pm}(n, q, \alpha)$  are

$$C_{\pm}(n, q, \alpha) = \frac{e^{i\frac{\pi}{2}\alpha} e^{\pm i\frac{\pi}{2}q} \Gamma\left(\frac{n-\alpha}{2}\right)}{2^{\alpha} \pi^{\frac{n}{2}} \Gamma\left(\frac{\alpha}{2}\right)}. \quad (\text{V}, 1; 10)$$

It is useful to register the particular case of the Theorem 4 corresponding to  $\alpha = 2k$ . We observe that the distributions  $H_{2k}(P \pm i0, n)$  are (causal, anticausal) elementary solutions of the homogeneous ultrahyperbolic operator iterated  $k$ -times (cf. [1], p. 41, formulae (II,3;13) and (II,4;4)).

**Theorem 5.** *The elementary solutions  $H_{2k}(P \pm i0, n)$  admit the integral representation*

$$H_{2k}(P \pm i0, n) = C_{\pm}(n, k, q) \gamma^{-1}\left(\frac{n-2k}{2}, P \pm i0, n\right) \lim_{\varepsilon \rightarrow 0} \int_0^{\infty} e^{-(P \pm i\varepsilon |x|^2)t} e^{\left(\frac{2k-n}{2}\right)e^{-t}} dt, \quad (\text{V}, 1; 11)$$

where the constants  $C_{\pm}(n, k, q)$  are

$$C_{\pm}(n, k, q) = \frac{e^{i\pi k} e^{\pm i\frac{\pi}{2}q} \Gamma\left(\frac{n-2k}{2}\right)}{2^{\alpha} \pi^{\frac{n}{2}} \Gamma(k)}. \quad (\text{V}, 1; 12)$$

**Note.** We know that  $H_0(P \pm i0, n) = \delta$  (cf. form. (II,1;8), p. 323, [7]).

Therefore, the thesis of Theorem 5, permit us, by taking limits in both members of (V,1;9) for  $k \rightarrow 0$ , obtain an integral representation of the  $\delta$ -measure of Dirac by means of an inverse Laplace transform.

## References.

- [1] Trione, S.E., *Distributional Products*, Cursos de Matemática, **3**, IAM - CONICET, Buenos Aires, Argentina, 1980.
- [2] Gelfand, I.M. and G.E. Shilov, *Generalized Functions, Vol.I*, Academic Press, New York, 1964.
- [3] Watson, G.N., *A treatise on the theory of Bessel functions*, Second Edition, Cambridge, University Press, 1947.



- [4] Bateman, *Manuscript project, Tables of Integral Transforms*, Vol.I, McGraw-Hill, New York, 1954.
- [5] Bogoliubov, N.N. and O.S. Parasiuk, *Über die multiplication der Kausal - functionen in der Quanten-theorie der Felder*, Acta Mathematica, **97**, 227-266, 1957.
- [6] Riesz, M., *L'integrale de Riemann-Liouville et le problème de Cauchy*, Acta Mathematica, **81**, pp. 1-223, 1953.
- [7] Trione, S.E., *On the Fourier Transform of Causal Distributions*, Studies in Applied Mathematics, **55**, 315-326, 1976.